Computational photonics — scattering problems in guided wave optics

Manfred Hammer*

Applied Analysis & Mathematical Physics
MESA+ Institute for Nanotechnology
University of Twente, The Netherlands

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* *Department of Applied Mathematics, University of Twente P.O. Box 217, 7500 AE Enschede, The Netherlands*
Phone: +31/53/489-3448 Fax: +31/53/489-4833 E-mail: m.hammer@math.utwente.nl

Photonics

"Photonics is the science and technology of generating, controlling, and detecting photons, particularly in the visible light and near infra-red spectrum.

… quantum optics, … optoelectronics …

… the overlap between all these fields and optics is unclear, and different definitions are used in different parts of the world and in different industries.

… goal of establishing an electronics of photons instead of electrons.

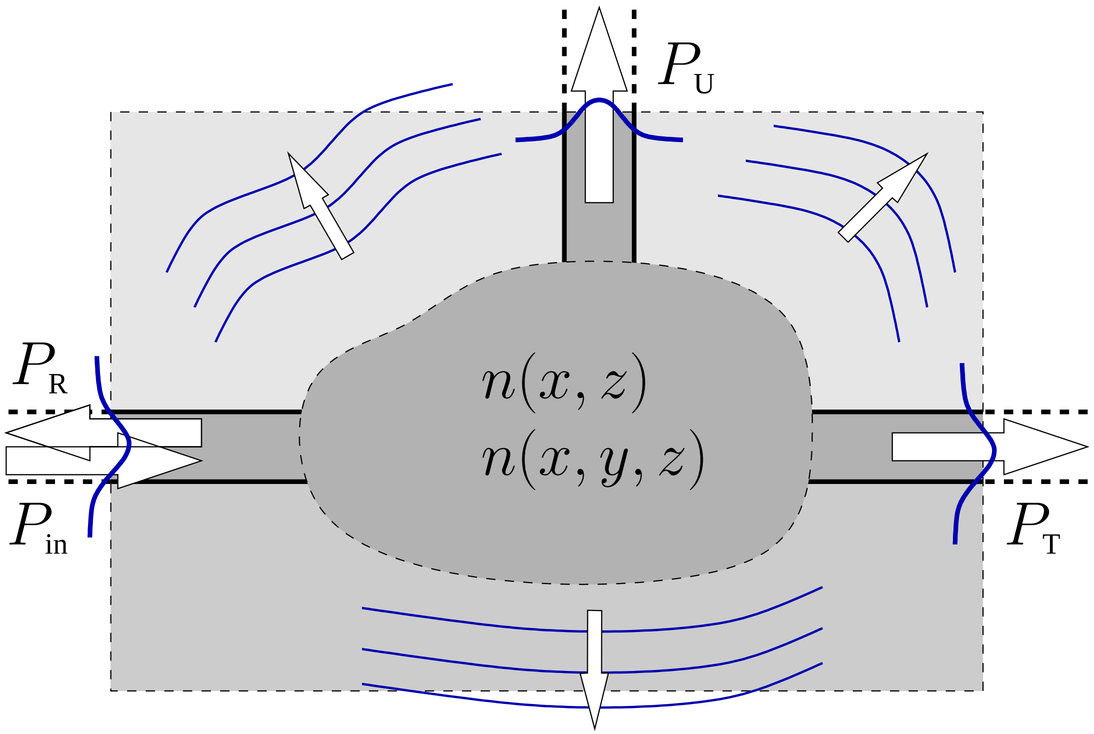
… strong interest in optical communication … usually based on laser light."

http://www.wikipedia.org

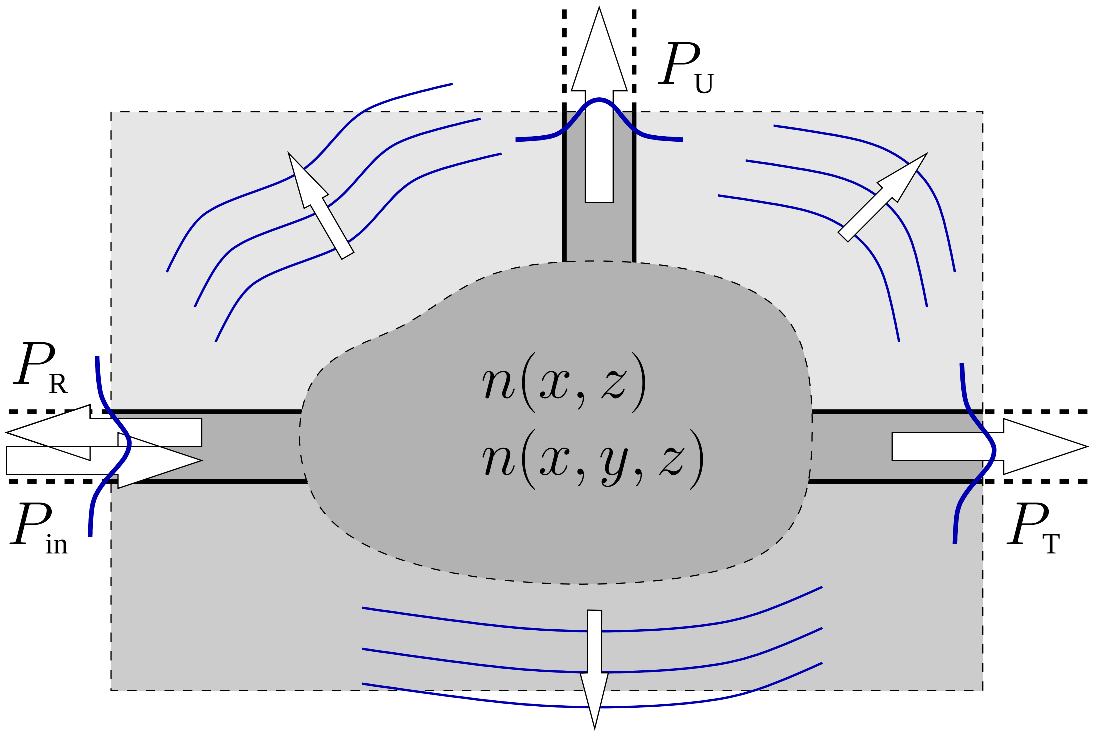
Outline

- Simulations in guided wave optics
  - Macroscopic Maxwell equations
  - Stationary and time-domain problems
  - 2-D problems
  - Modes of dielectric waveguides
- Quadridirectional eigenmode propagation
  - Problem setting
  - Eigenmode expansion
  - Algebraic procedure
- 2-D PSTM model
  - Probing evanescent fields
  - Hole defect in a slab waveguide
  - Waveguide Bragg grating: experiment & 2-D model
  - Resonant defect cavity
- Optical switching by NEMS-actuated resonator arrays

Abstract scattering problem



Abstract scattering problem



Typical parameters:

- vacuum wavelength
  $\lambda \in [400, 700]$ nm (visible light),
  $\lambda \approx 1.3\,\mu$m, $1.55\,\mu$m
  (optical fibers, attenuation min.),
- refractive indices $n \in [1, 3.4]$,
  small attenuation
  (transparent dielectrica).

- Interesting domain: $(10\,\lambda — 100\,\lambda)^d$, $d = 2, 3$ (2-D, 3-D).
- Details: $\approx \lambda/10$, $\approx \lambda/100$.
- Influx and outflux: guided & nonguided waves $\leftrightsquigarrow$ boundary conditions.

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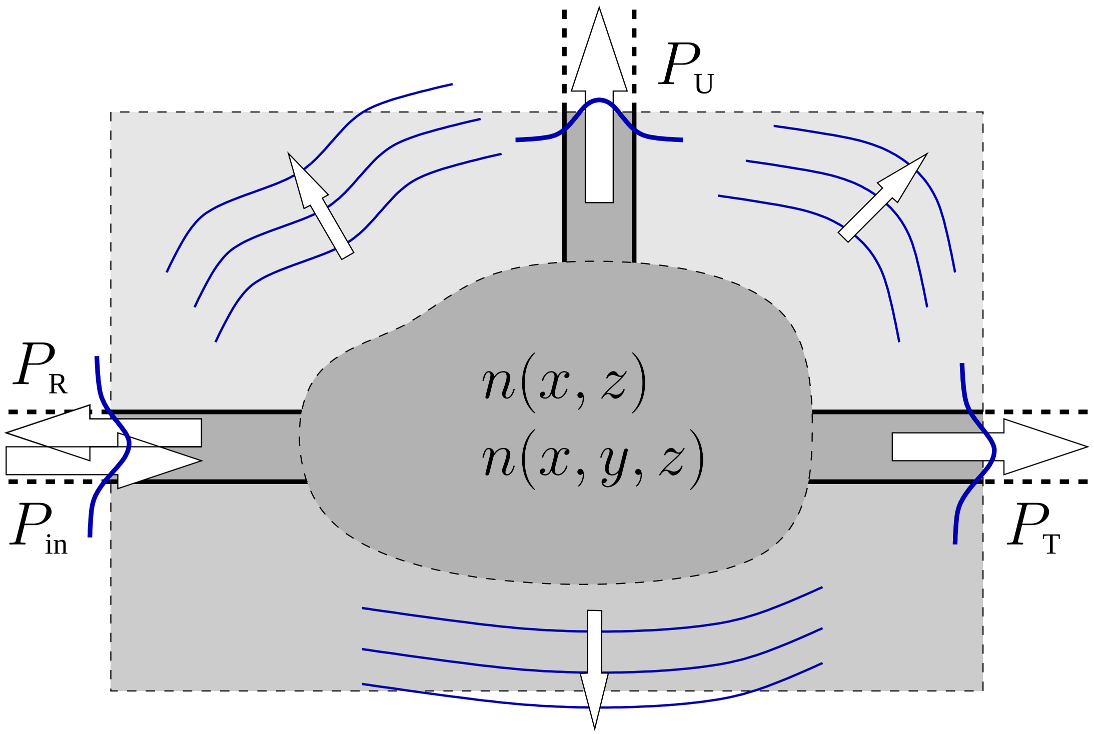

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- Influx and outflux: guided & nonguided waves $\longleftrightarrow$ boundary conditions.

- Emphasis: device concepts, design.

Macroscopic Maxwell equations

… for the electromagnetic fields

$$\boldsymbol{\mathcal{E}}(x,y,z,t) = \frac{1}{2}\left(\boldsymbol{E}(x,y,z)\,\mathrm{e}^{\mathrm{i}\omega t} + \boldsymbol{E}^*(x,y,z)\,\mathrm{e}^{-\mathrm{i}\omega t}\right),$$

$\hat{=}\ \boldsymbol{\mathcal{B}}, \boldsymbol{B},\ \boldsymbol{\mathcal{D}}, \boldsymbol{D},\ \boldsymbol{\mathcal{H}}, \boldsymbol{H},\ \boldsymbol{\mathcal{P}}, \boldsymbol{P},\ \boldsymbol{\mathcal{M}}, \boldsymbol{M},\ \omega = k\mathrm{c} = 2\pi\mathrm{c}/\lambda,$

in frequency domain form, SI (source-free):

$$\mathrm{curl}\,\boldsymbol{E} = -\mathrm{i}\omega\boldsymbol{B}, \quad \mathrm{curl}\,\boldsymbol{H} = \mathrm{i}\omega\boldsymbol{D}, \quad \mathrm{div}\,\boldsymbol{D} = 0, \quad \mathrm{div}\,\boldsymbol{B} = 0,$$

$$\boldsymbol{B} = \mu_0(\boldsymbol{H} + \boldsymbol{M}), \quad \boldsymbol{D} = \epsilon_0\boldsymbol{E} + \boldsymbol{P}.$$

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$\hat{=}$ $\boldsymbol{\mathcal{B}}$, $\boldsymbol{B}$, $\boldsymbol{\mathcal{D}}$, $\boldsymbol{D}$, $\boldsymbol{\mathcal{H}}$, $\boldsymbol{H}$, $\boldsymbol{\mathcal{P}}$, $\boldsymbol{P}$, $\boldsymbol{\mathcal{M}}$, $\boldsymbol{M}$, $\omega = k\mathrm{c} = 2\pi\mathrm{c}/\lambda$,

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$$\boldsymbol{B} = \mu_0(\boldsymbol{H} + \boldsymbol{M}), \quad \boldsymbol{D} = \epsilon_0\boldsymbol{E} + \boldsymbol{P}.$$

Typical media:

- uncharged, no free currents and charges,
- nonmagnetic at optical frequencies, $\boldsymbol{M} = 0$,
- dielectrica, susceptibilities $\hat{\chi}^{(j)}(x,y,z;\omega)$:

$$P_j = \epsilon_0\Big(\sum_k \chi^{(1)}_{j,k}E_k + \sum_{k,l}\chi^{(2)}_{j,k,l}E_kE_l + \sum_{k,l,m}\chi^{(3)}_{j,k,l,m}E_kE_lE_m \ldots\Big).$$

Convention: eliminate $\boldsymbol{D}$ and $\boldsymbol{B}$ $\rightsquigarrow$ formulation in $\boldsymbol{E}$ and $\boldsymbol{H}$.

Linear problems

$$\boldsymbol{P} = \epsilon_0 \hat{\chi}^{(1)} \boldsymbol{E}, \quad \boldsymbol{D} = \epsilon_0 (1 + \hat{\chi}^{(1)}) \boldsymbol{E} = \epsilon_0 \hat{\epsilon} \boldsymbol{E}\,;$$

$\hat{\epsilon} = 1 + \hat{\chi}^{(1)}$: relative permittivity.

Simplest case: isotropic, lossless dielectrica; refractive index n :

$$\hat{\epsilon} = \epsilon 1, \quad \epsilon = n^2, \quad n(x, y, z; \omega) \in \mathbb{R}.$$

Frequently n is piecewise constant $\leftrightsquigarrow$ interface conditions for $\boldsymbol{E}$, $\boldsymbol{H}$.

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Frequently n is piecewise constant $\leftrightsquigarrow$ interface conditions for $\boldsymbol{E}$, $\boldsymbol{H}$.

Complications:

- anisotropic media, $\hat{\epsilon} \not\sim 1$ (crystals, ordered materials),
- attenuation, $\hat{\epsilon}^\dagger \neq \hat{\epsilon}$, $n \notin \mathbb{R}$,
- nonlinear problems, $\hat{\chi}^{(2)}$, $\hat{\chi}^{(3)}$ $\ldots$

Stationary problems

Continuous wave excitation, ω a given parameter:

$$\operatorname{curl} \boldsymbol{E} = -\mathrm{i}\omega\mu_0 \boldsymbol{H}, \quad \operatorname{curl} \boldsymbol{H} = \mathrm{i}\omega\epsilon_0\hat{\epsilon}\boldsymbol{E}, \quad \operatorname{div} \hat{\epsilon}\boldsymbol{E} = 0, \quad \operatorname{div} \boldsymbol{H} = 0,$$

or

$$\operatorname{curl}\operatorname{curl} \boldsymbol{E} = k^2\hat{\epsilon}\boldsymbol{E}, \quad \operatorname{curl} \hat{\epsilon}^{-1}\operatorname{curl} \boldsymbol{H} = k^2\boldsymbol{H},$$

$\boldsymbol{E}(x, y, z) \in \mathbb{C}^3$, $\boldsymbol{H}(x, y, z) \in \mathbb{C}^3$.

Helmholtz solver:

Given $\hat{\epsilon}$ and an optical “influx”, find $\boldsymbol{E}$, $\boldsymbol{H}$ on a computational domain, subject to suitable boundary conditions.

Scans over ω $\rightsquigarrow$ Spectral data.

Time dependent modeling

Propagation of time dependent signals, pulsed excitation:

$$\operatorname{curl}\boldsymbol{\mathcal{E}} = -\mu_0\partial_t\boldsymbol{\mathcal{H}}, \quad \operatorname{curl}\boldsymbol{\mathcal{H}} = \epsilon_0\hat{\epsilon}\partial_t\boldsymbol{\mathcal{E}}, \quad \operatorname{div}\hat{\epsilon}\boldsymbol{\mathcal{E}} = 0, \quad \operatorname{div}\boldsymbol{\mathcal{H}} = 0,$$

or

$$\operatorname{curl}\operatorname{curl}\boldsymbol{\mathcal{E}} = -\frac{1}{c^2}\hat{\epsilon}\partial_t^2\boldsymbol{\mathcal{E}}, \quad \operatorname{curl}\hat{\epsilon}^{-1}\operatorname{curl}\boldsymbol{\mathcal{H}} = -\frac{1}{c^2}\partial_t^2\boldsymbol{\mathcal{H}},$$

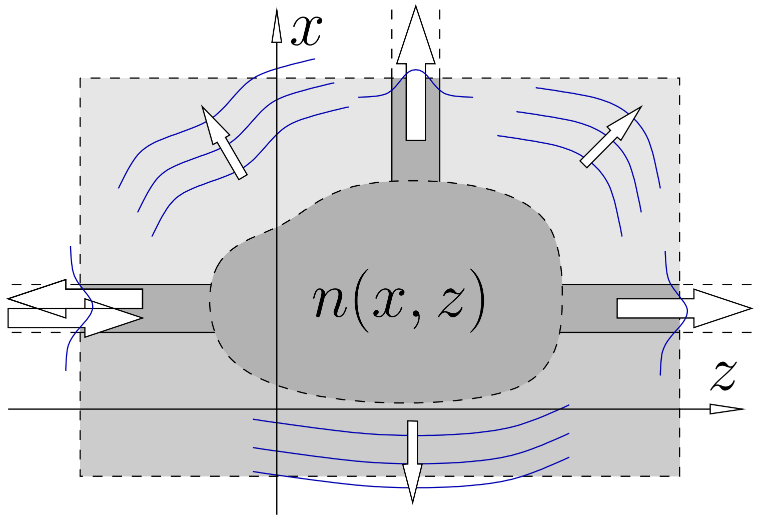
$\boldsymbol{\mathcal{E}}(x,y,z,t) \in \mathbb{R}^3$, $\boldsymbol{\mathcal{H}}(x,y,z,t) \in \mathbb{R}^3$.

Time domain solver:

Given $\hat{\epsilon}$ and an optical “influx” signal, find $\boldsymbol{\mathcal{E}}$, $\boldsymbol{\mathcal{H}}$ on a computational domain within a certain time interval, subject to suitable boundary conditions.

Fourier transform with respect to time $\rightsquigarrow$ spectral data.

2-D problems



$\partial_y \epsilon = 0,\ \epsilon(x,z) = n^2(x,z),$
$\partial_y \boldsymbol{E} = 0,\ \partial_y \boldsymbol{H} = 0;$
equations split into two subsets:

TE, E_y, H_x, and H_z, principal component E_y:

$$\mathrm{i}\omega\mu_0 H_x = \partial_z E_y, \quad \mathrm{i}\omega\mu_0 H_z = -\partial_x E_y, \quad \mathrm{i}\omega\epsilon_0\epsilon E_y = \partial_z H_x - \partial_x H_z,$$

or

$$\partial_x^2 E_y, + \partial_z^2 E_y + k^2 \epsilon E_y = 0.$$

TM, H_y, E_x, and E_z, principal component H_y:

$$\mathrm{i}\omega\epsilon_0\epsilon E_x = -\partial_z H_y, \quad \mathrm{i}\omega\epsilon_0\epsilon E_z = \partial_x H_y, \quad -\mathrm{i}\omega\mu_0 H_y = \partial_z E_x - \partial_x E_z,$$

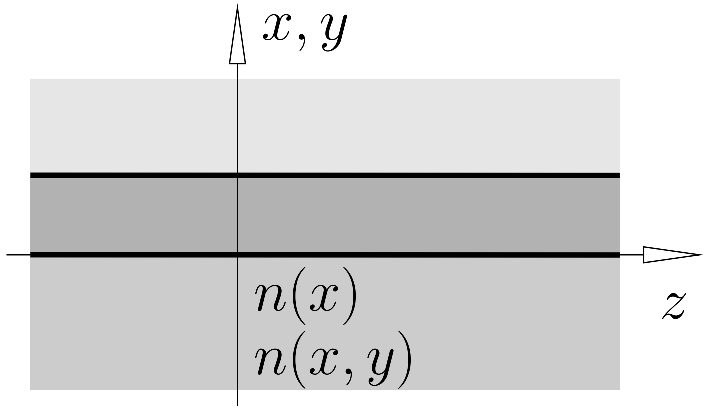
or

$$\partial_x \frac{1}{\epsilon} \partial_x H_y, + \partial_z \frac{1}{\epsilon} \partial_z H_y + k^2 H_y = 0.$$

$\partial_z \epsilon = 0$, $\partial_z n = 0$, modal solutions with profile $\tilde{\boldsymbol{E}}, \tilde{\boldsymbol{H}}$, and propagation constant β;

$\boldsymbol{E}(x,y,z) = \tilde{\boldsymbol{E}}(x,y)\,\mathrm{e}^{-\mathrm{i}\beta z}$,

$\boldsymbol{H}(x,y,z) = \tilde{\boldsymbol{H}}(x,y)\,\mathrm{e}^{-\mathrm{i}\beta z}$:

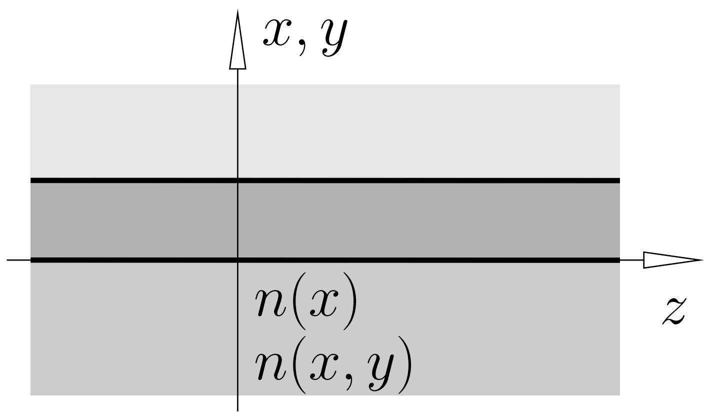


$$\boldsymbol{\mathcal{E}}(x,y,z,t) = \mathrm{Re}\,\tilde{\boldsymbol{E}}(x,y)\,\mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z}, \quad \boldsymbol{\mathcal{H}}(x,y,z,t) = \mathrm{Re}\,\tilde{\boldsymbol{H}}(x,y)\,\mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z}.$$

Guided modes: $\displaystyle\iint |\tilde{\boldsymbol{E}}|^2 \mathrm{d}x\,\mathrm{d}y < \infty, \quad \iint |\tilde{\boldsymbol{H}}|^2 \mathrm{d}x\,\mathrm{d}y < \infty.$

Mode solver: Eigenvalue problem for $\tilde{\boldsymbol{E}}$, $\tilde{\boldsymbol{H}}$, and β.

Dielectric waveguides



$\partial_z \epsilon = 0$, $\partial_z n = 0$, modal solutions with profile $\tilde{\boldsymbol{E}}, \tilde{\boldsymbol{H}}$, and propagation constant β;

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$$\boldsymbol{H}(x, y, z) = \tilde{\boldsymbol{H}}(x, y)\, \mathrm{e}^{-\mathrm{i}\beta z}.$$

$$\boldsymbol{\mathcal{E}}(x, y, z, t) = \mathrm{Re}\, \tilde{\boldsymbol{E}}(x, y)\, \mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z}, \quad \boldsymbol{\mathcal{H}}(x, y, z, t) = \mathrm{Re}\, \tilde{\boldsymbol{H}}(x, y)\, \mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z}.$$

Guided modes: $\iint |\tilde{\boldsymbol{E}}|^2 \mathrm{d}x\, \mathrm{d}y < \infty, \quad \iint |\tilde{\boldsymbol{H}}|^2 \mathrm{d}x\, \mathrm{d}y < \infty.$

Mode solver: Eigenvalue problem for $\tilde{\boldsymbol{E}}$, $\tilde{\boldsymbol{H}}$, and β.

- Basis for all kinds of design considerations,
- basis fields for various types of perturbational simulations,
- input / output fields for Helmholtz & time domain solvers.

Guided modes, 2-D, TE

$$\left(\partial_x^2 + k^2 n^2(x)\right) \tilde{E}_y = \beta^2 \tilde{E}_y\,, \quad \int |\tilde{E}_y|^2 \mathrm{d}x < \infty\,, \quad \tilde{H}_x = \frac{-\beta}{\omega\mu_0}\tilde{E}_y\,, \quad \tilde{H}_z = \frac{\mathrm{i}}{\omega\mu_0}\partial_x \tilde{E}_y\,,$$

$$\boldsymbol{\mathcal{E}}(x,z,t) = \mathrm{Re}\begin{pmatrix} 0 \\ \tilde{E}_y \\ 0 \end{pmatrix}(x)\,\mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z}, \qquad \boldsymbol{\mathcal{H}}(x,z,t) = \mathrm{Re}\begin{pmatrix} \tilde{H}_x \\ 0 \\ \tilde{H}_z \end{pmatrix}(x)\,\mathrm{e}^{\mathrm{i}\omega t - \mathrm{i}\beta z},$$

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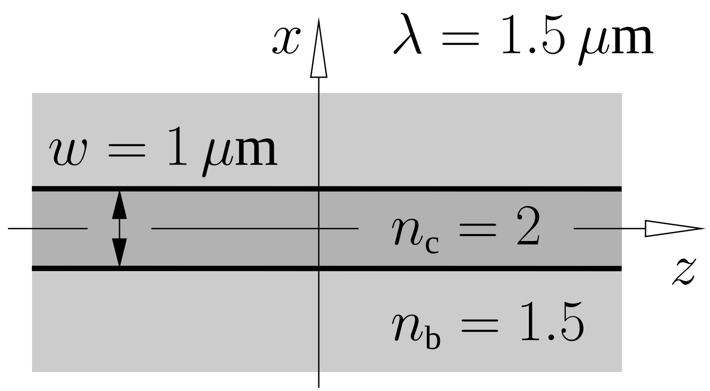
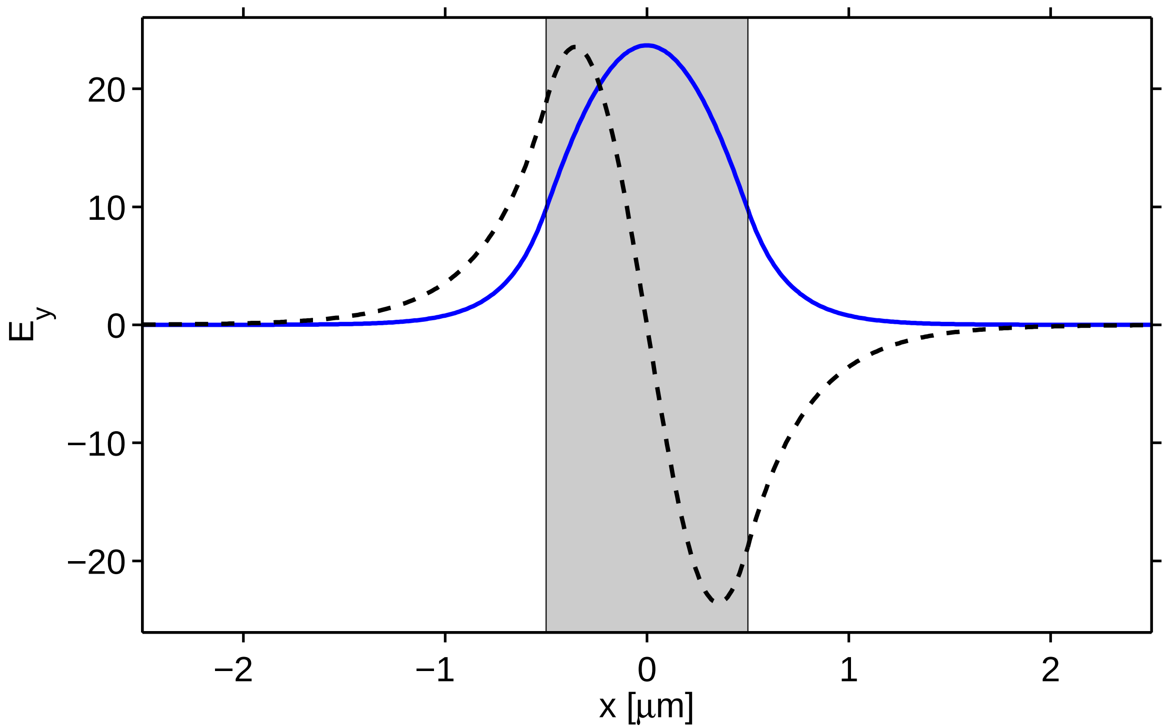
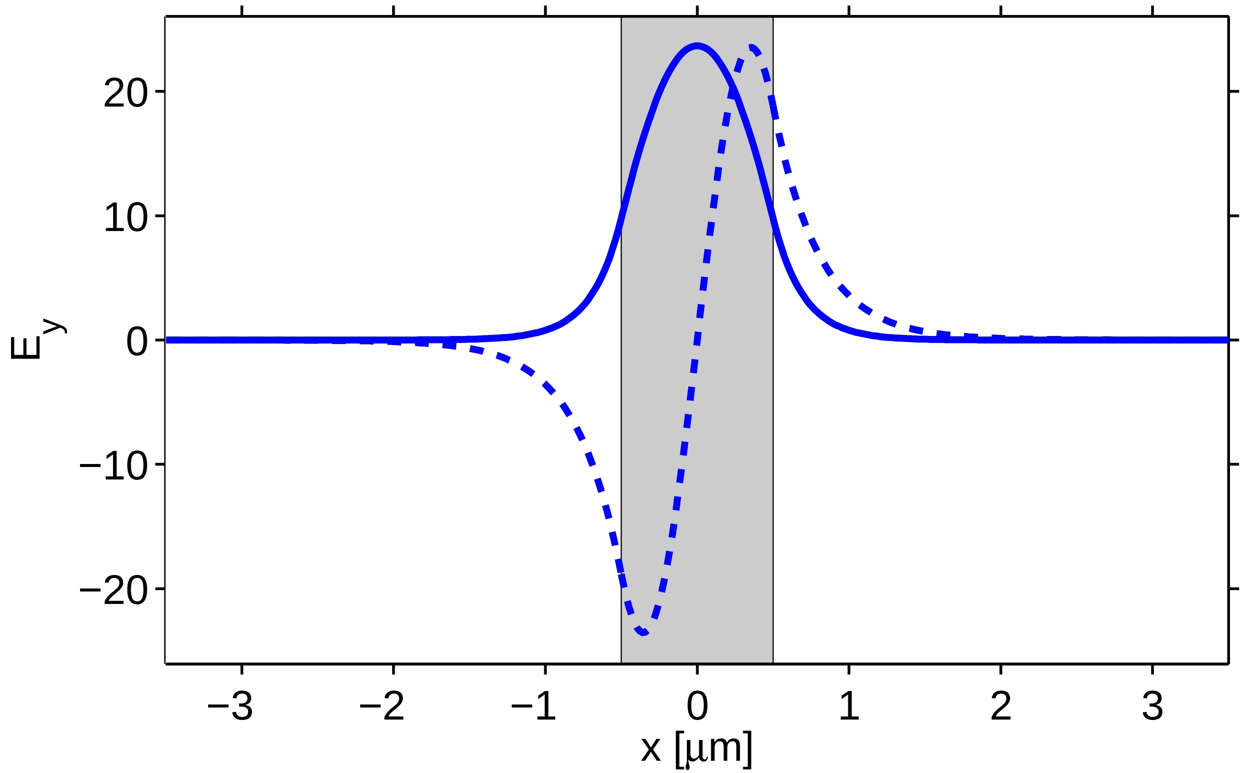


$\beta_0/k = 1.923$, $\beta_1/k = 1.697$.



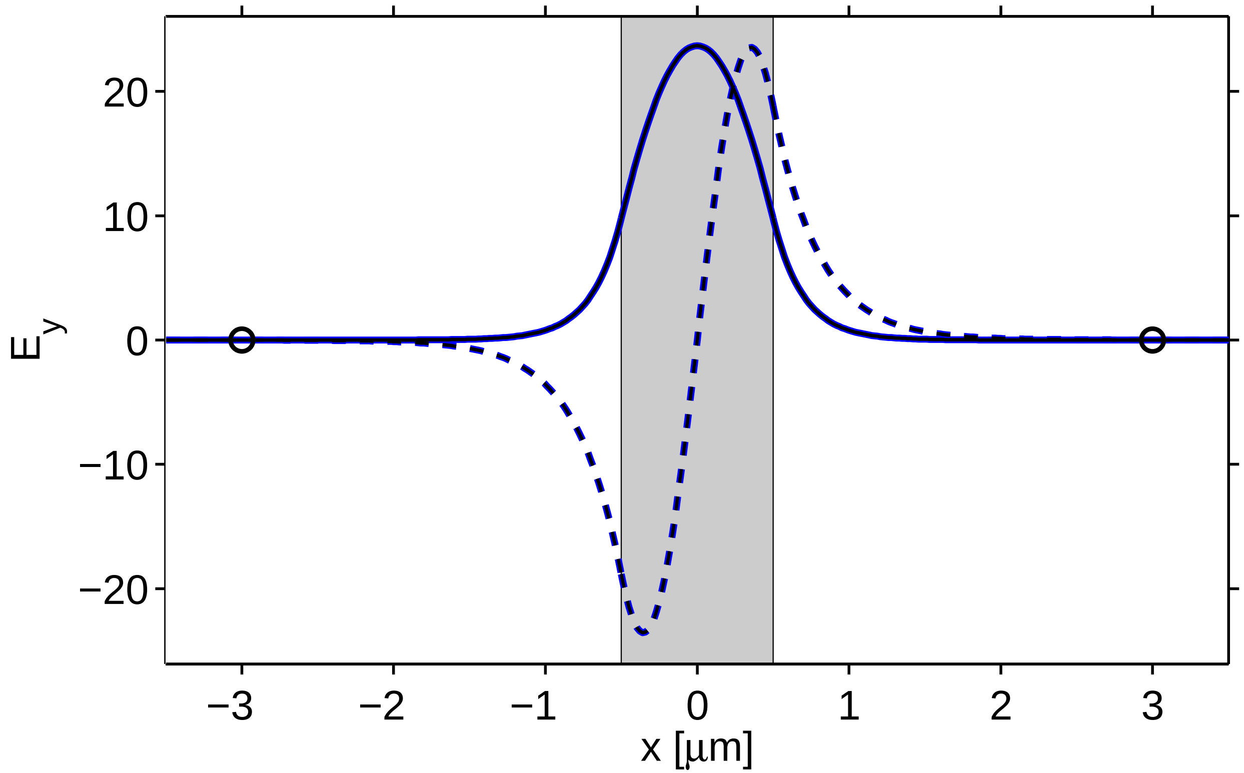
- Mode propagation
- “Coupling” of modes

Modal basis sets



$$\left(\partial_x^2 + k^2 n^2(x)\right) \tilde{E}_y = \beta^2 \tilde{E}_y\,, \qquad \int |\tilde{E}_y|^2 \mathrm{d}x < \infty.$$

Modal basis sets



$$\left(\partial_x^2 + k^2 n^2(x)\right) \tilde{E}_y = \beta^2 \tilde{E}_y \,, \qquad \int_{x_0}^{x_1} |\tilde{E}_y|^2 \mathrm{d}x = 1 \,, \qquad \tilde{E}_y(x_0) = \tilde{E}_y(x_1) = 0 \,.$$

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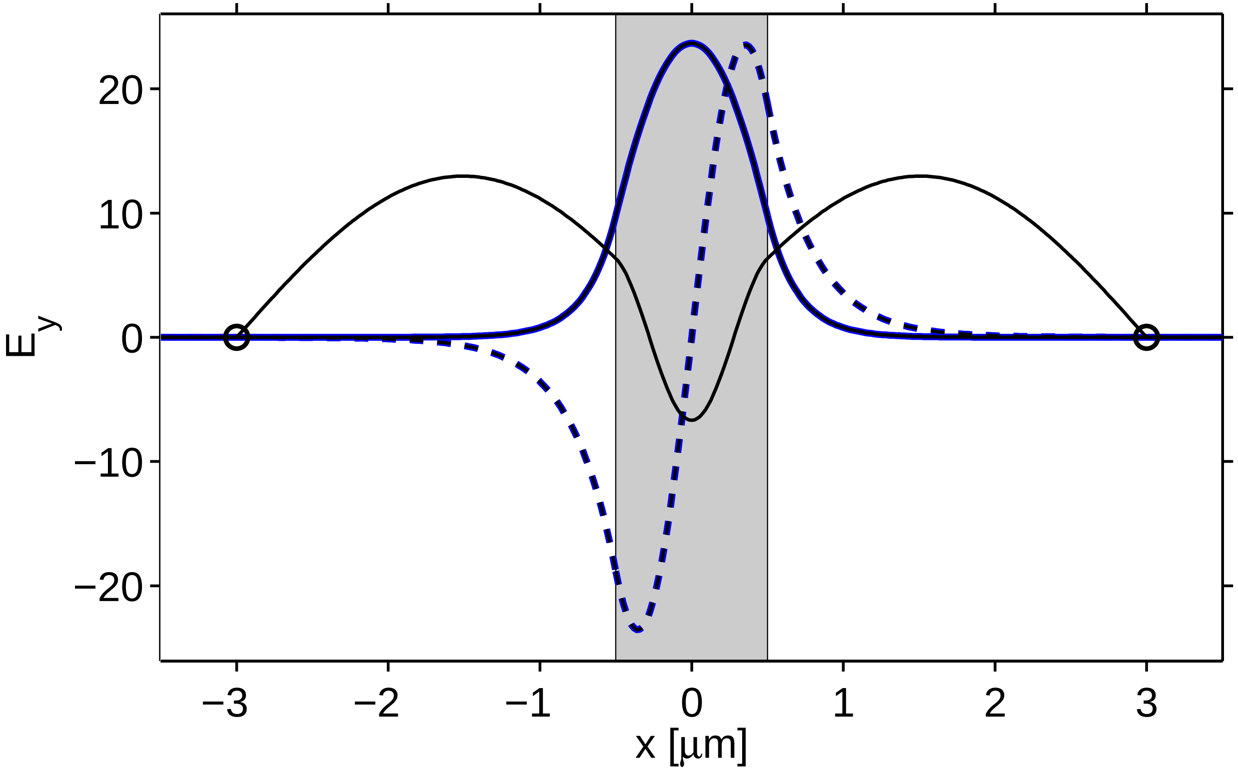

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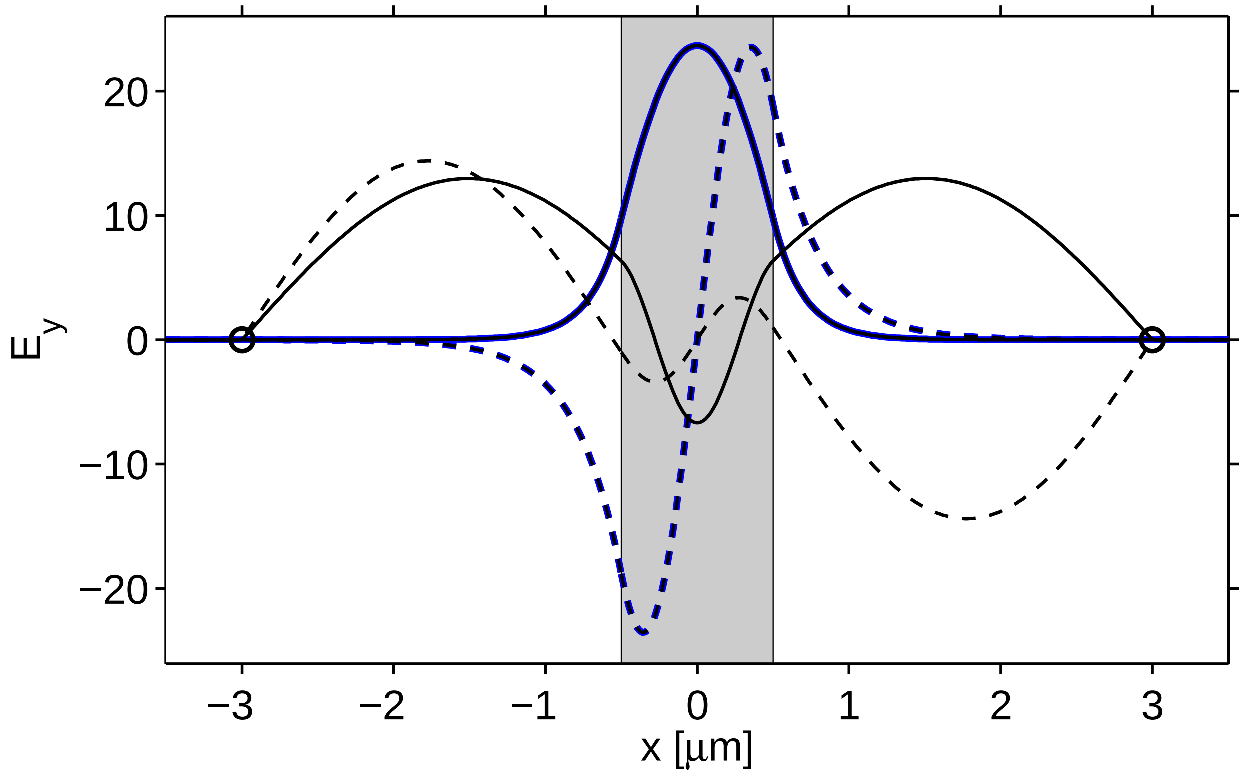

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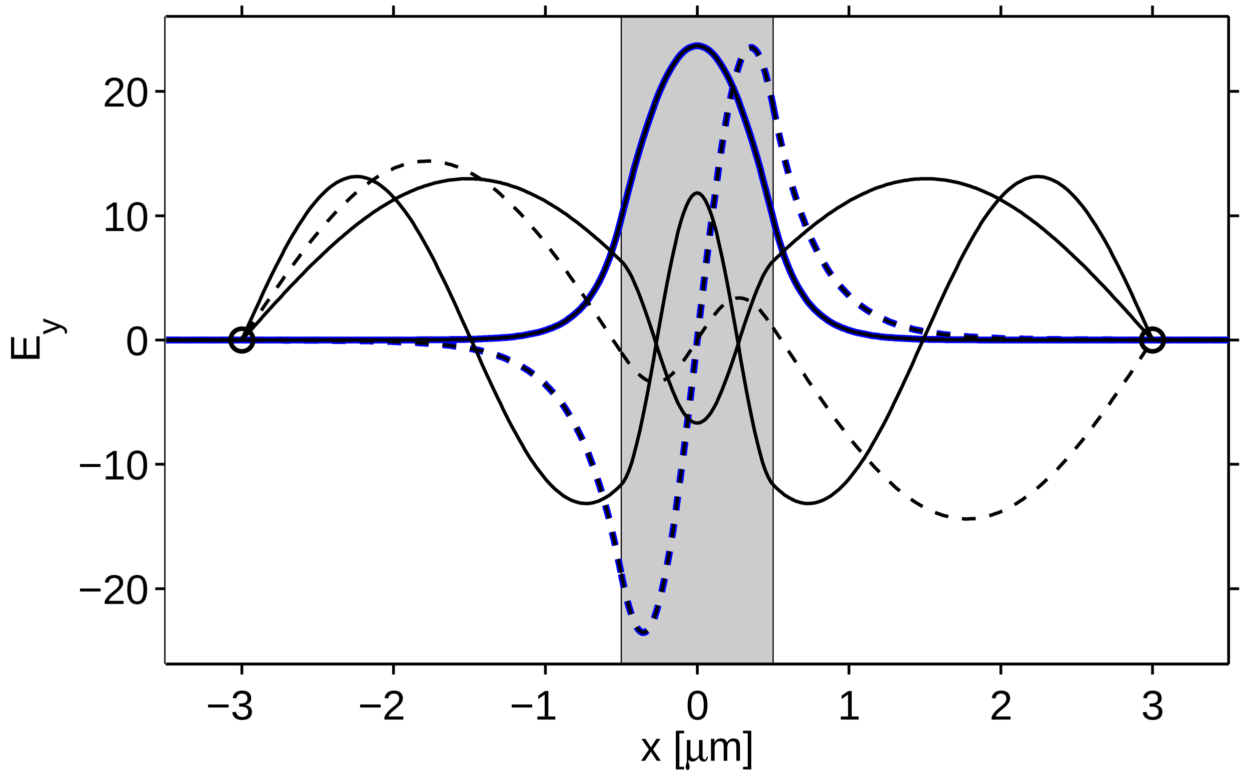

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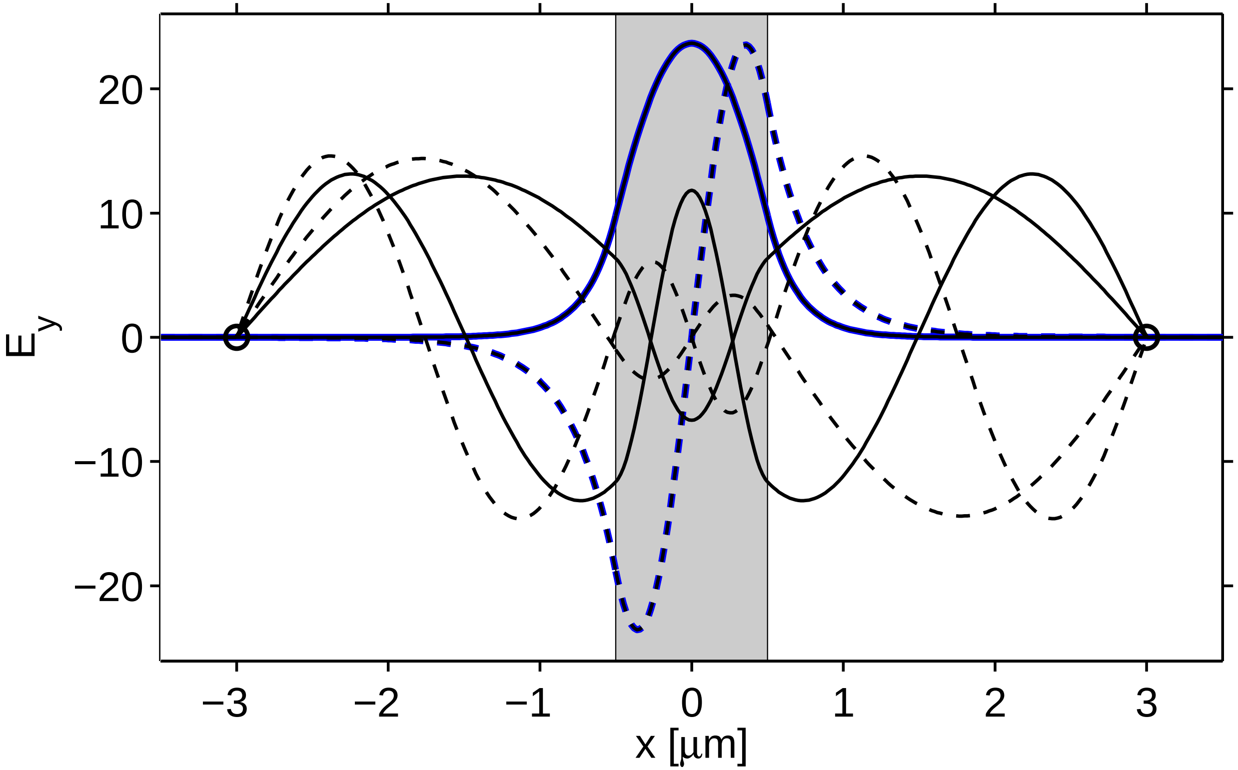

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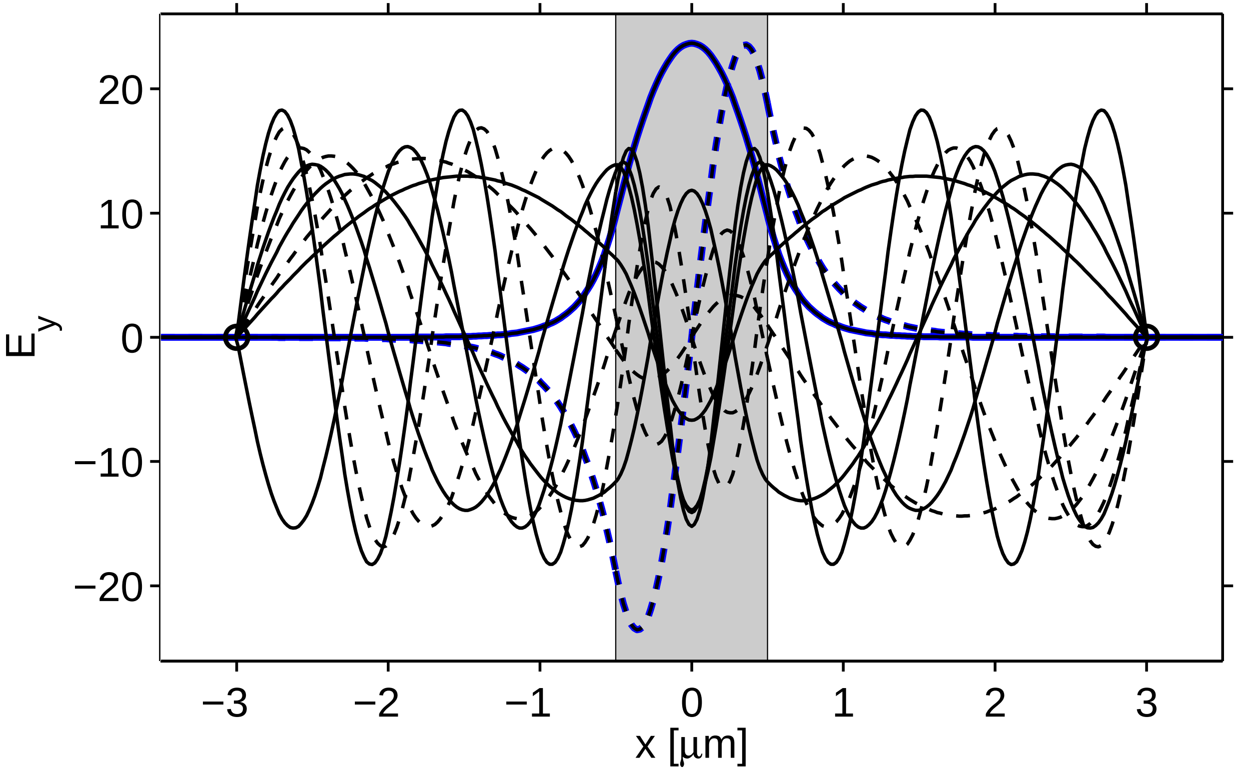

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$\{\beta_j, E_{y,j}\}:$ complete discrete set of eigenmodes, expansion basis.

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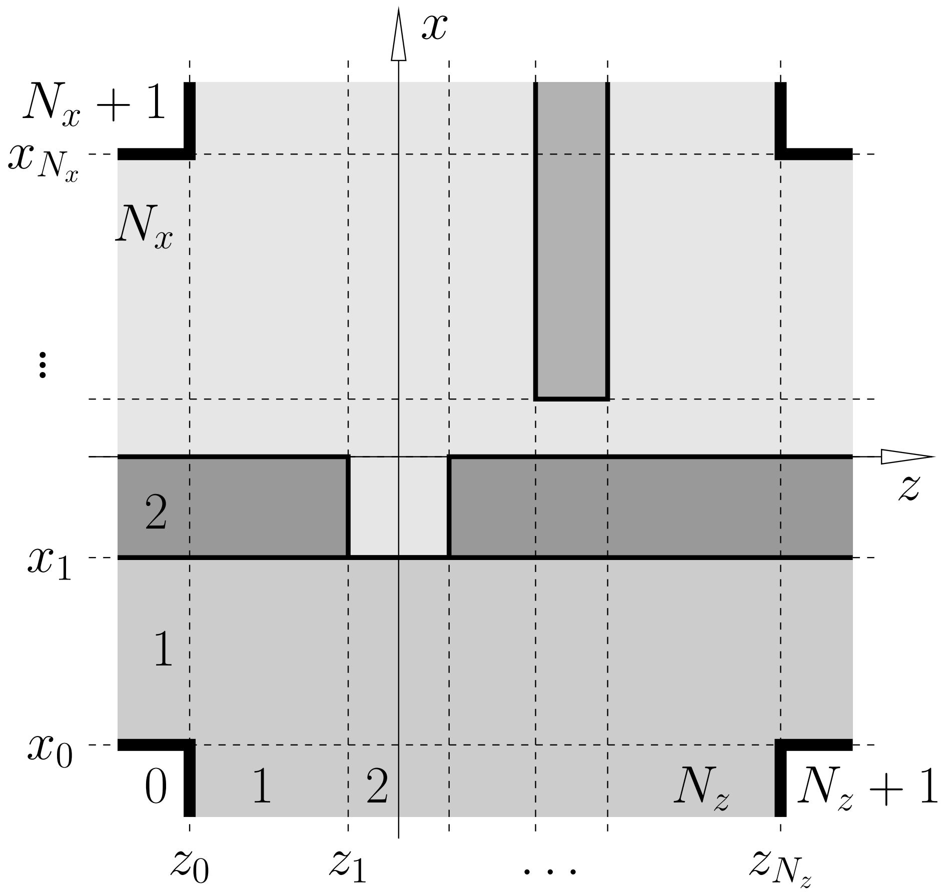
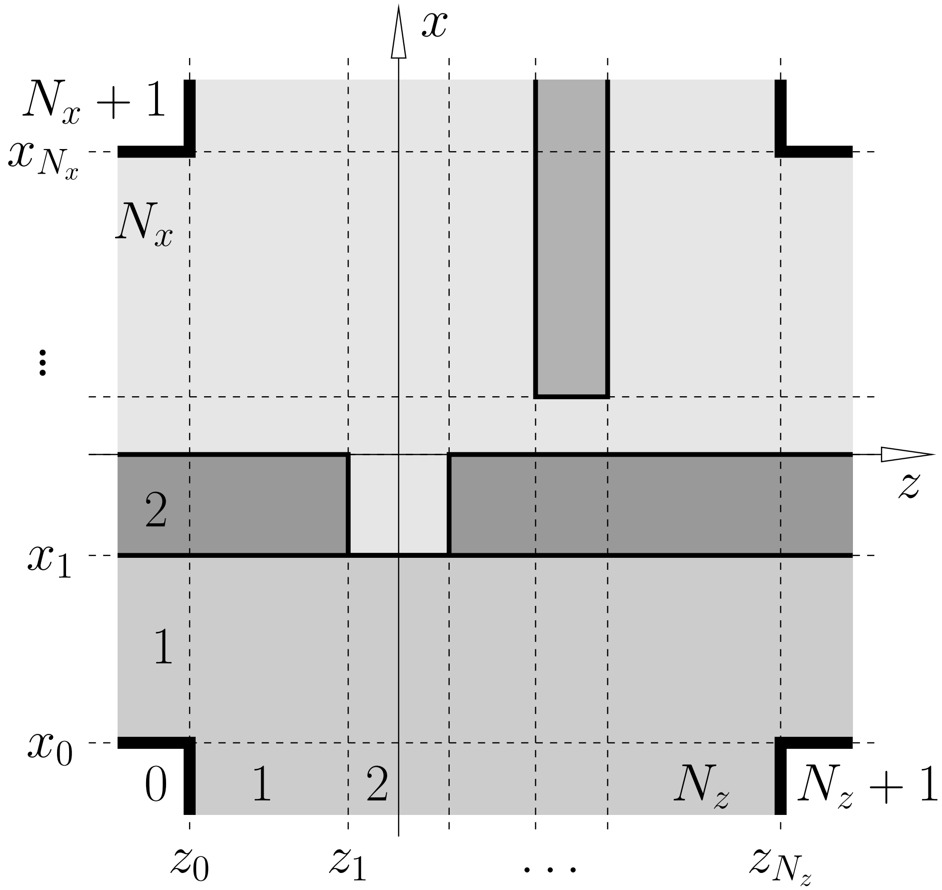


- 2-D TE/TM Helmholtz problem, vacuum wavelength $\lambda = 2\pi/k$.

QUEP simulations, problem setting



- 2-D TE/TM Helmholtz problem, vacuum wavelength $\lambda = 2\pi/k$.
- Piecewise constant, rectangular refractive index distribution; linear, lossless materials.

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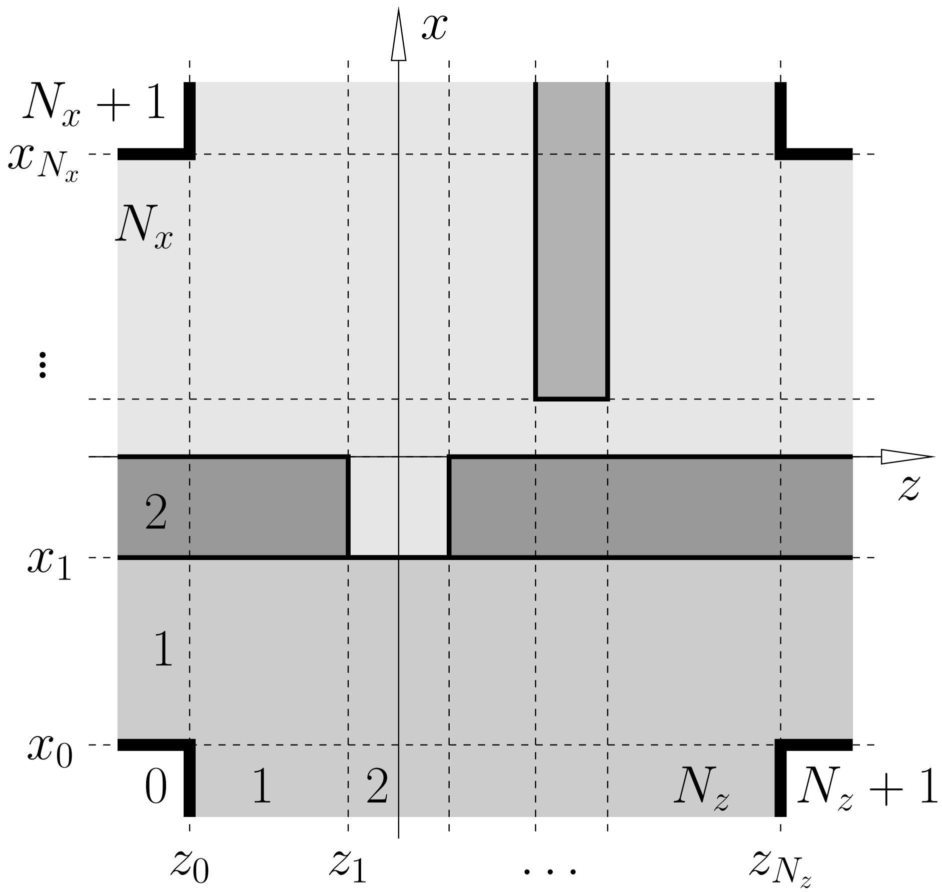

- 2-D TE/TM Helmholtz problem, vacuum wavelength $\lambda = 2\pi/k$.
- Piecewise constant, rectangular refractive index distribution; linear, lossless materials.
- Rectangular interior computational domain, influx & outflux across all four boundaries, outwards homogeneous external regions.

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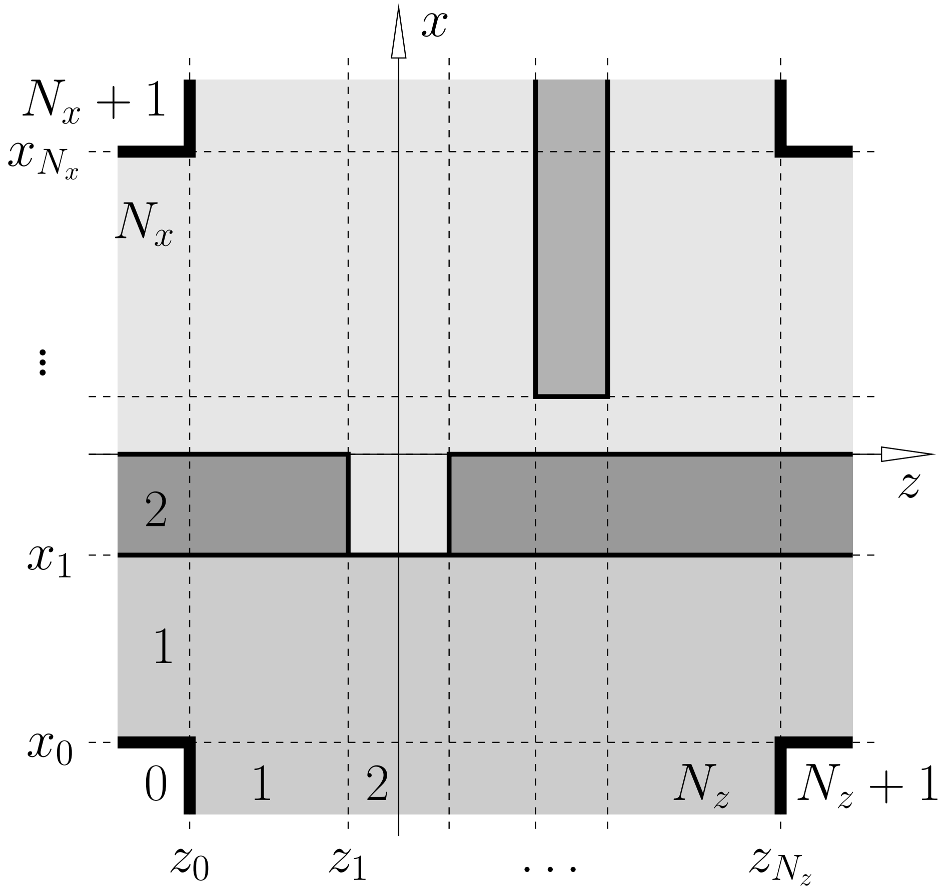
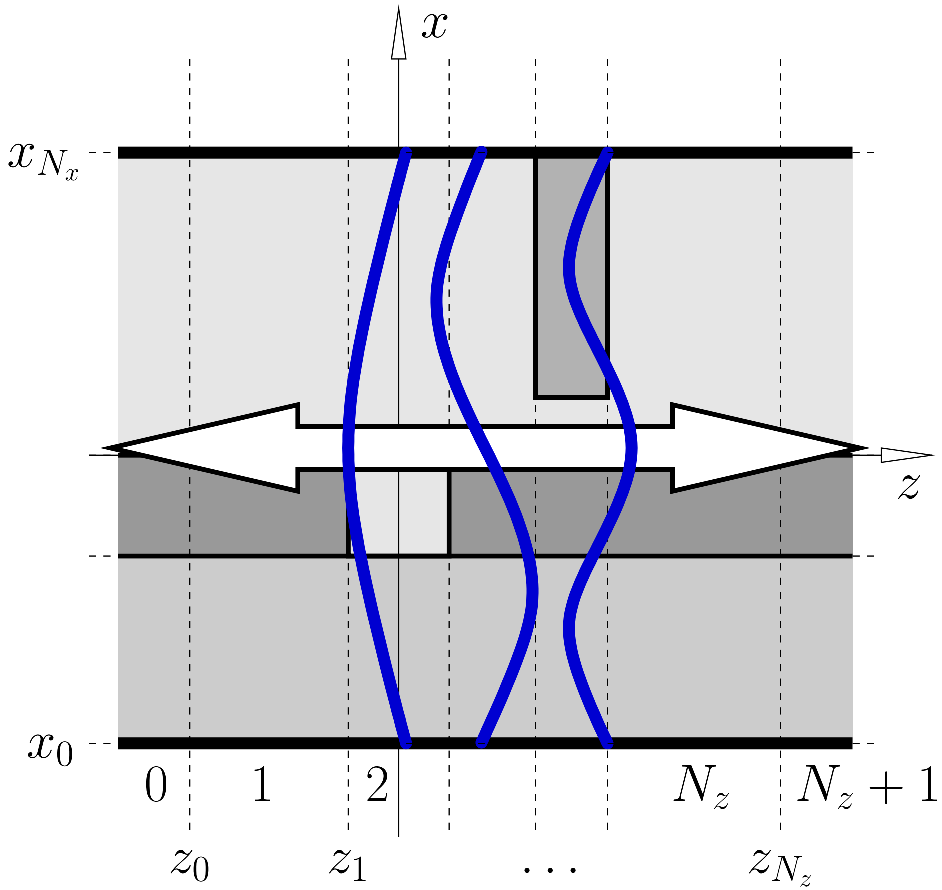


- 2-D TE/TM Helmholtz problem, vacuum wavelength $\lambda = 2\pi/k$.
- Piecewise constant, rectangular refractive index distribution; linear, lossless materials.
- Rectangular interior computational domain, influx & outflux across all four boundaries, outwards homogeneous external regions.
- Assumption $E_y = 0$, $H_y = 0$ on the corner points and on the external border lines is reasonable for the problems under investigation.

Ingredients: “horizontal” modes

Basis fields,
defined by Dirichlet boundary conditions $E_y = 0$ (TE) or $H_y = 0$ (TM):



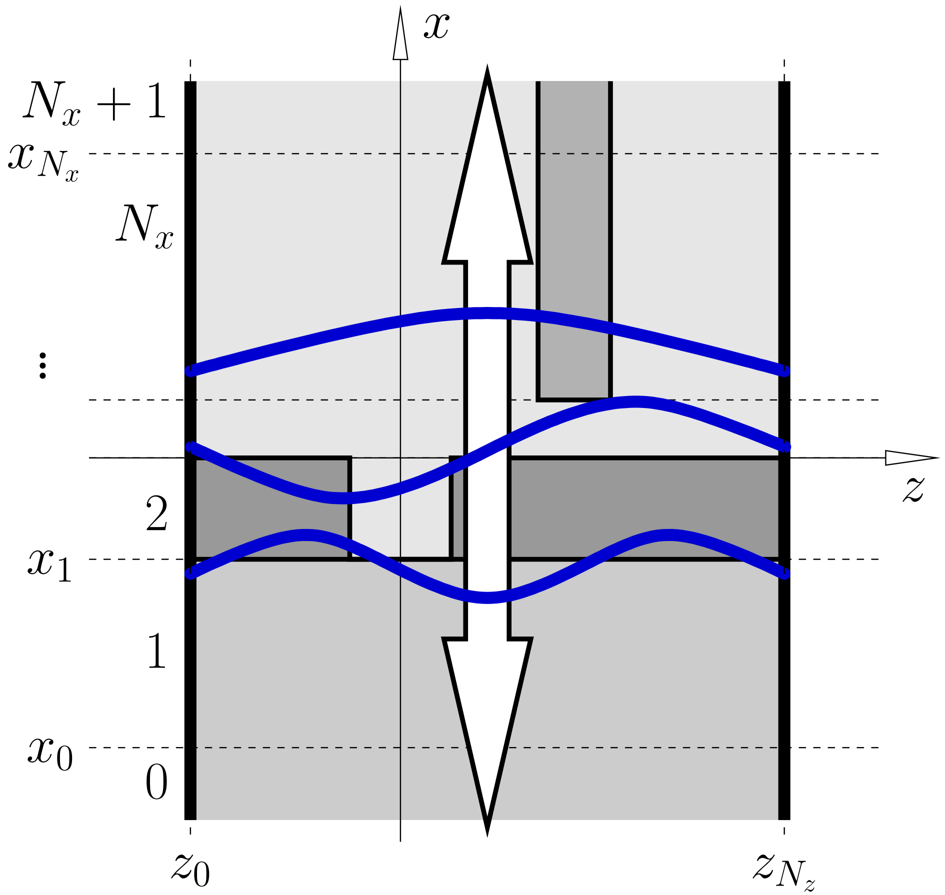
Horizontally traveling eigenmodes:

M_x profiles	$\boldsymbol{\psi}^d_{s,m}(x)$
and propagation constants	$\pm\beta_{s,m}$

of order m, on slice s,
for propagation directions $d =$ f, b,

Ingredients: "vertical" modes

Basis fields,
defined by Dirichlet boundary conditions $E_y = 0$ (TE) or $H_y = 0$ (TM):

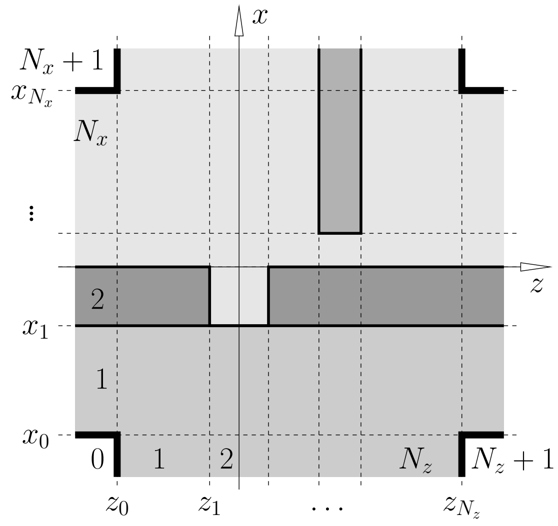


. . . *and* vertically traveling fields:

M_z profiles	$\hat{\psi}^d_{l,m}(z)$
and propagation constants	$\pm\hat{\beta}_{l,m}$

of order m, on layer l,
for propagation directions $d =$ u, d.

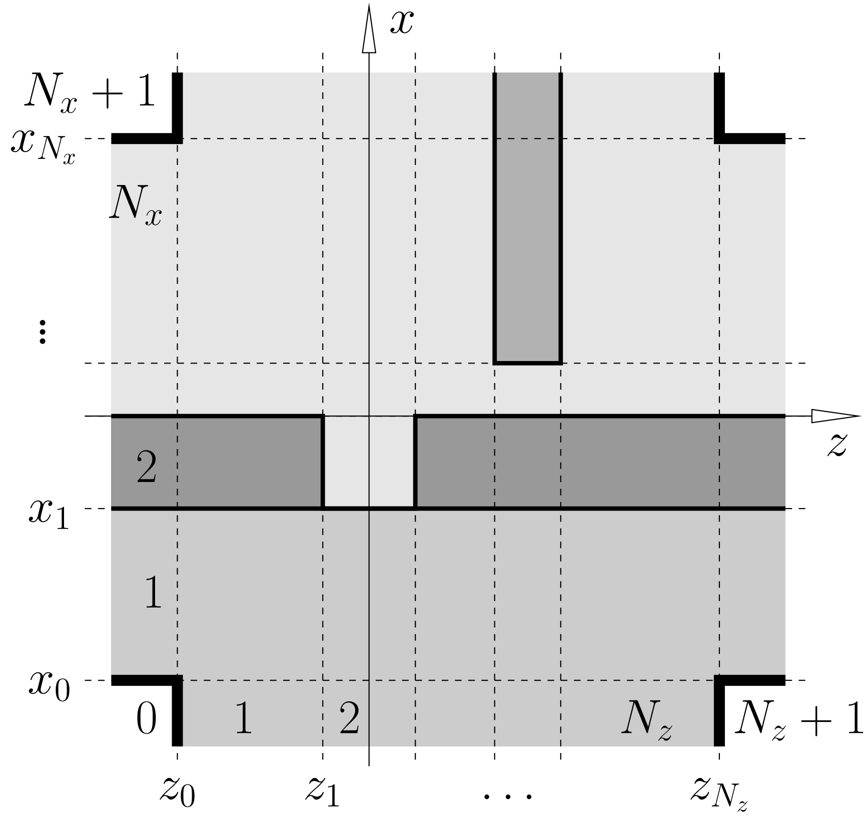
Eigenmode expansion



Ansatz for the optical field,
for $z_{s-1} \le z \le z_s$, $s = 1, \dots, N_z$,
and $x_{l-1} \le x \le x_l$, $l = 1, \dots, N_x$:

$$\begin{pmatrix} \boldsymbol{\mathcal{E}} \\ \boldsymbol{\mathcal{H}} \end{pmatrix}(x,z,t) = \mathrm{Re}\left\{ \sum_{m=0}^{M_x-1} F_{s,m} \boldsymbol{\psi}^{\mathrm{f}}_{s,m}(x)\, \mathrm{e}^{-\mathrm{i}\beta_{s,m}(z-z_{s-1})} + \sum_{m=0}^{M_x-1} B_{s,m} \boldsymbol{\psi}^{\mathrm{b}}_{s,m}(x)\, \mathrm{e}^{+\mathrm{i}\beta_{s,m}(z-z_s)} + \sum_{m=0}^{M_z-1} U_{l,m} \hat{\boldsymbol{\psi}}^{\mathrm{u}}_{l,m}(z)\, \mathrm{e}^{-\mathrm{i}\hat{\beta}_{l,m}(x-x_{l-1})} + \sum_{m=0}^{M_z-1} D_{l,m} \hat{\boldsymbol{\psi}}^{\mathrm{d}}_{l,m}(z)\, \mathrm{e}^{+\mathrm{i}\hat{\beta}_{l,m}(x-x_l)} \right\} \mathrm{e}^{\mathrm{i}\omega t}.$$

Eigenmode expansion

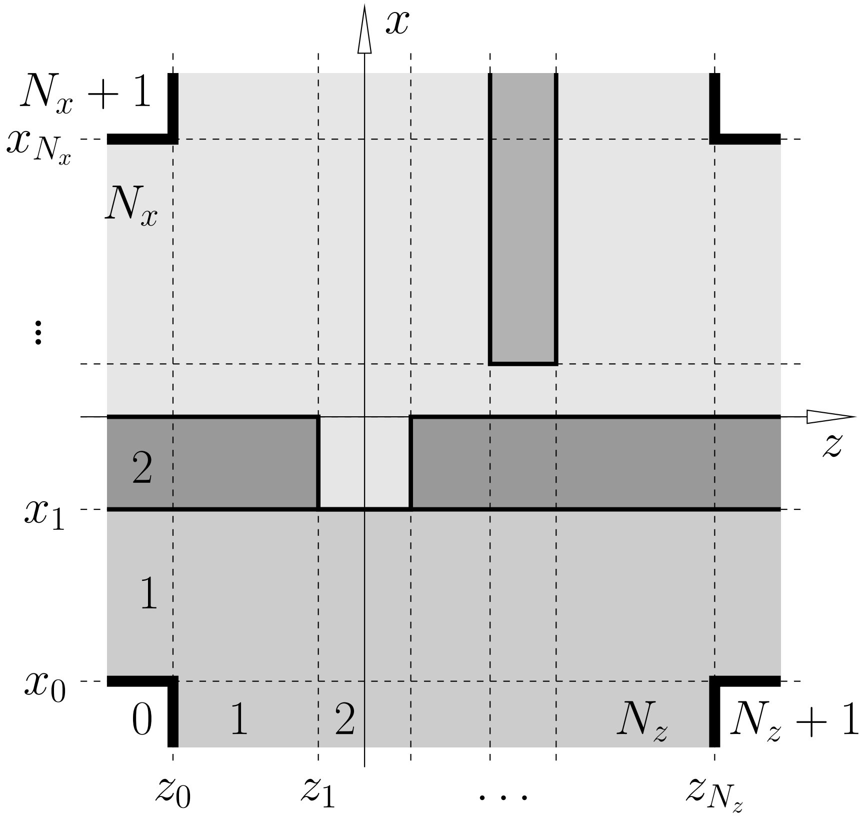


Mode products $\leftrightarrow$ normalization, projection:

$$(\boldsymbol{E}_1, \boldsymbol{H}_1; \boldsymbol{E}_2, \boldsymbol{H}_2) = \frac{1}{4} \int (E^*_{1,x} H_{2,y} - E^*_{1,y} H_{2,x} + H^*_{1,y} E_{2,x} - H^*_{1,x} E_{2,y})\, \mathrm{d}x \,,$$

$$\langle \boldsymbol{E}_1, \boldsymbol{H}_1; \boldsymbol{E}_2, \boldsymbol{H}_2 \rangle = \frac{1}{4} \int (E^*_{1,y} H_{2,z} - E^*_{1,z} H_{2,y} + H^*_{1,z} E_{2,y} - H^*_{1,y} E_{2,z})\, \mathrm{d}z \,.$$

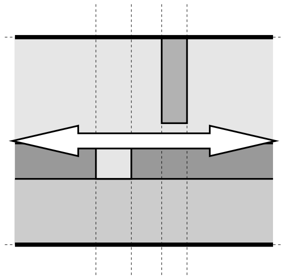
Algebraic procedure



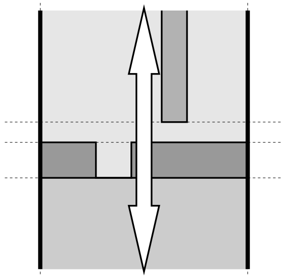
- Consistent bidirectional projection at all interfaces $\longrightarrow$ linear system of equations in $\{F_{s,m}, B_{s,m}, U_{l,m}, D_{l,m}\}$.
- Influx: $\boldsymbol{F}_0,\ \boldsymbol{B}_{N_x+1},\ \boldsymbol{U}_0,\ \boldsymbol{D}_{N_z+1} \ \longrightarrow$ RHS, given.
- Outflux: $\boldsymbol{B}_0,\ \boldsymbol{F}_{N_x+1},\ \boldsymbol{D}_0,\ \boldsymbol{U}_{N_z+1} \ \longrightarrow$ primary unknowns.

Algebraic procedure

- "Exact" mode profiles $\longrightarrow$ interior problems decouple:

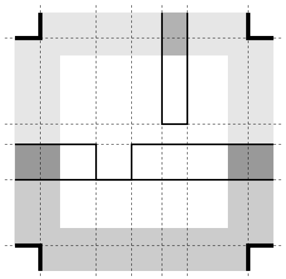


Solve for $\boldsymbol{F}_2, \ldots, \boldsymbol{F}_{N_z}$ and $\boldsymbol{B}_1, \ldots, \boldsymbol{B}_{N_z-1}$
in terms of $\boldsymbol{F}_1$ and $\boldsymbol{B}_{N_z}$ $\longrightarrow$ BEP I.



Solve for $\boldsymbol{U}_2, \ldots, \boldsymbol{U}_{N_x}$ and $\boldsymbol{D}_1, \ldots, \boldsymbol{D}_{N_x-1}$
in terms of $\boldsymbol{U}_1$ and $\boldsymbol{D}_{N_x}$ $\longrightarrow$ BEP II.

- Continuity of E and H on outer interfaces:



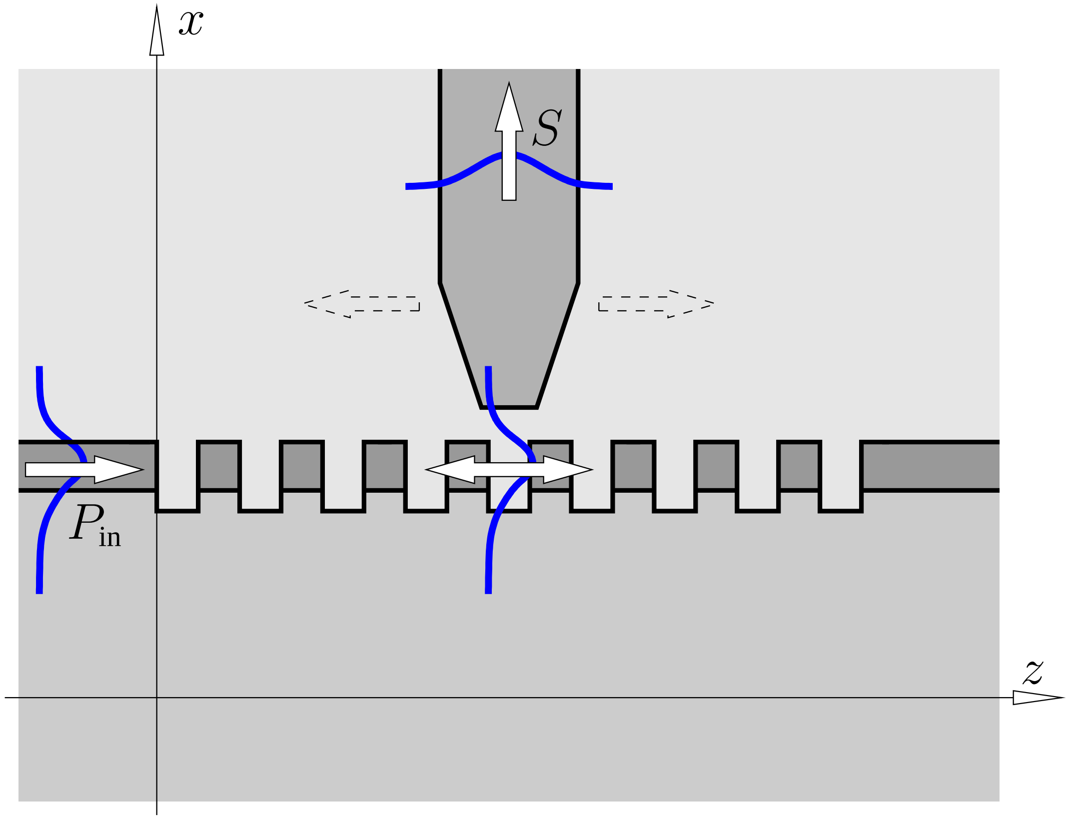
Interior BEP solutions
$+$ equations at $z = z_0, z_{N_z},\ x = x_0, x_{N_x}$
$\longrightarrow \boldsymbol{B}_0,\ \boldsymbol{F}_{N_x+1},\ \boldsymbol{D}_0,\ \boldsymbol{U}_{N_z+1}$.

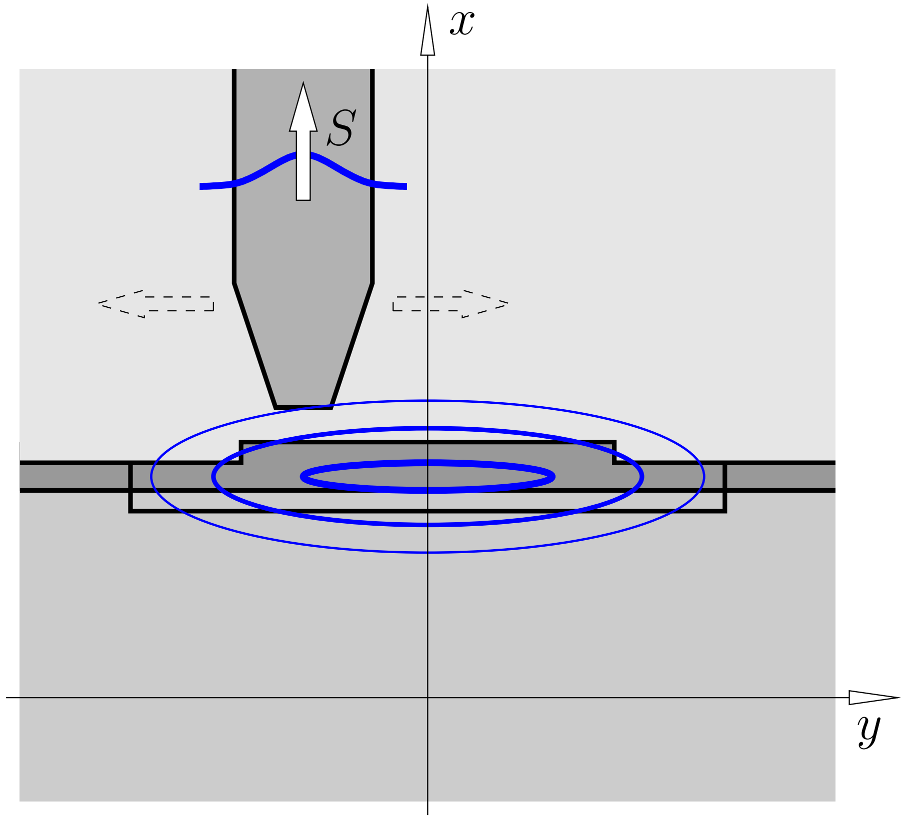
"QUadridirectional Eigenmode Propagation method" (QUEP).

Outline

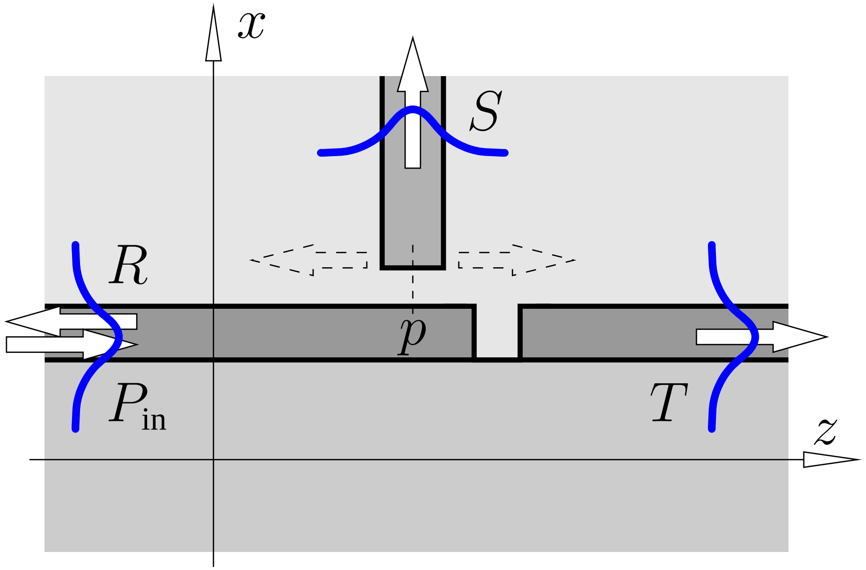
- Simulations in guided wave optics
  - Macroscopic Maxwell equations
  - Stationary and time-domain problems
  - 2-D problems
  - Modes of dielectric waveguides
- Quadridirectional eigenmode propagation
  - Problem setting
  - Eigenmode expansion
  - Algebraic procedure
- 2-D PSTM model
  - Probing evanescent fields
  - Hole defect in a slab waveguide
  - Waveguide Bragg grating: experiment & 2-D model
  - Resonant defect cavity
- Optical switching by NEMS-actuated resonator arrays

Photon scanning tunneling microscopy (PSTM) of photonic devices

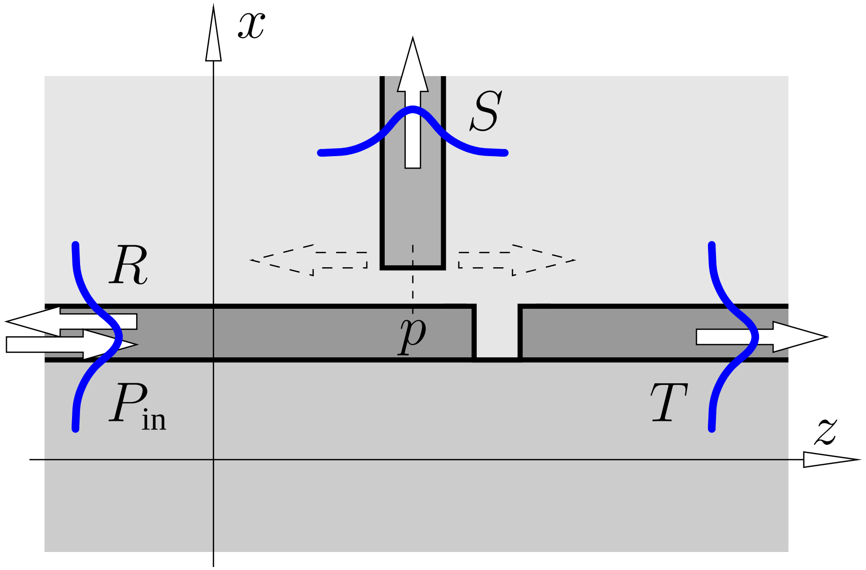




Simplified 2-D model



Simplified 2-D model

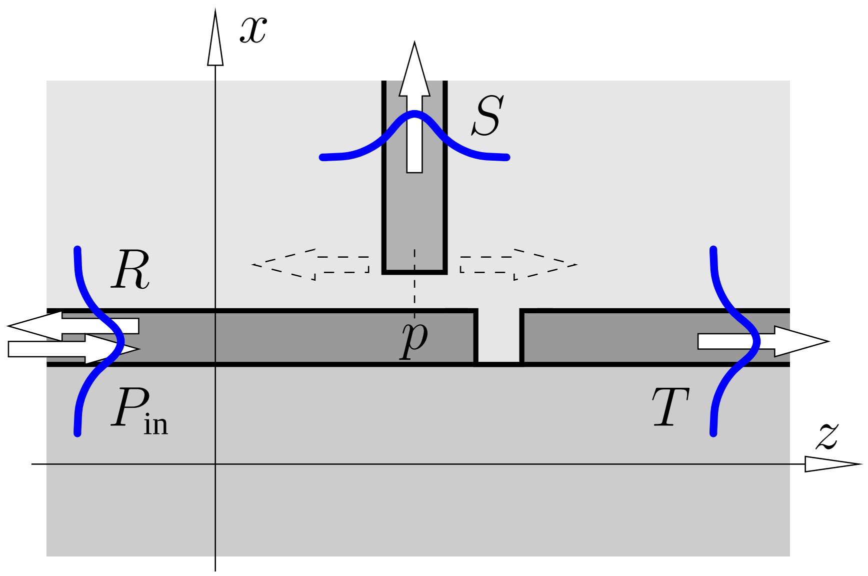


$P_{\text{in}} = 1. \quad S(p) = ?$

$R(p), T(p) = ?$

$S(p) \overset{?}{\longleftrightarrow} \text{(optical field)}\big|_p$

Simplified 2-D model



$P_{\text{in}} = 1. \quad S(p) = ?$

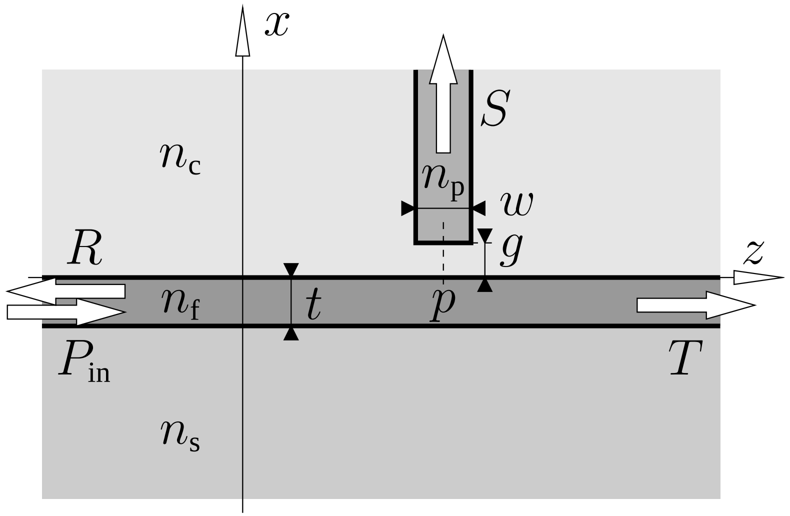
$R(p), T(p) = ?$

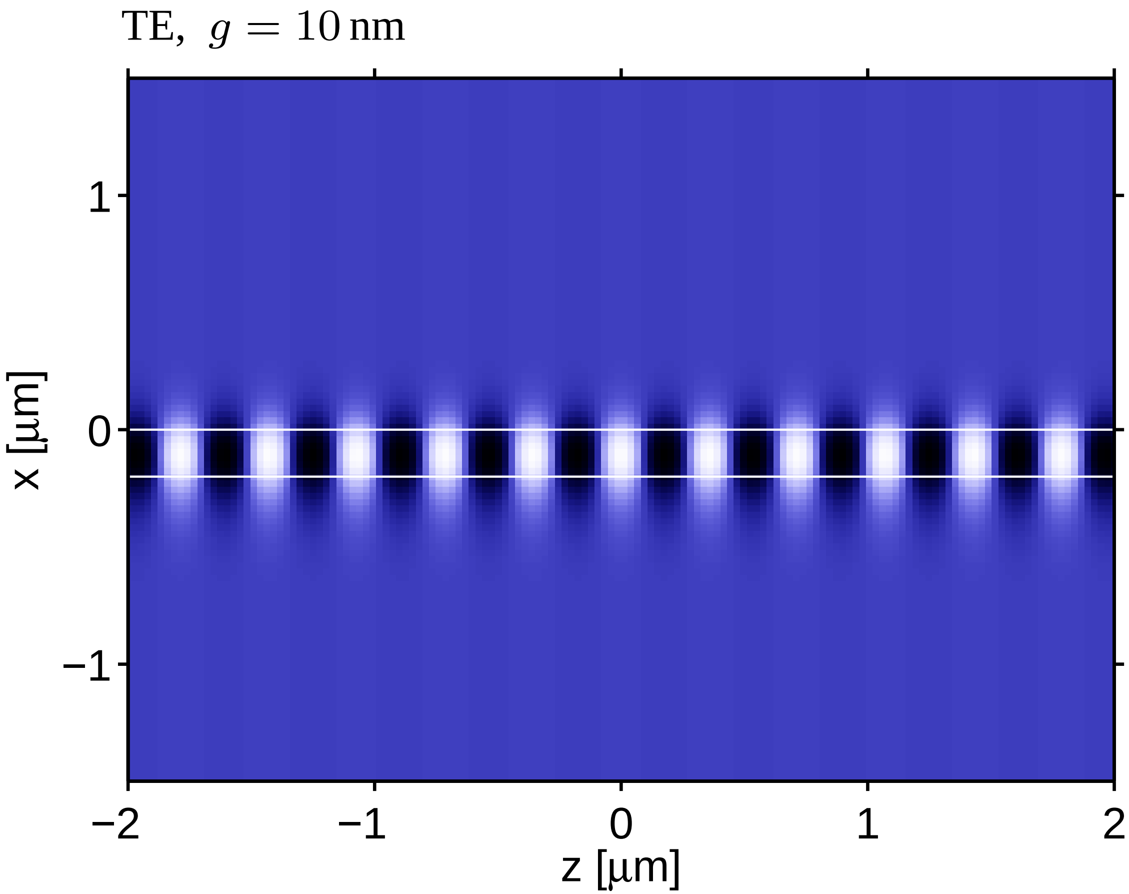
$S(p) \overset{?}{\longleftrightarrow} (\text{optical field})\big|_p$

Fixed optical frequency, rectangular piecewise constant refractive index,
horizontal & vertical guided wave in- and outflux

$\longrightarrow$ *QUEP simulations*

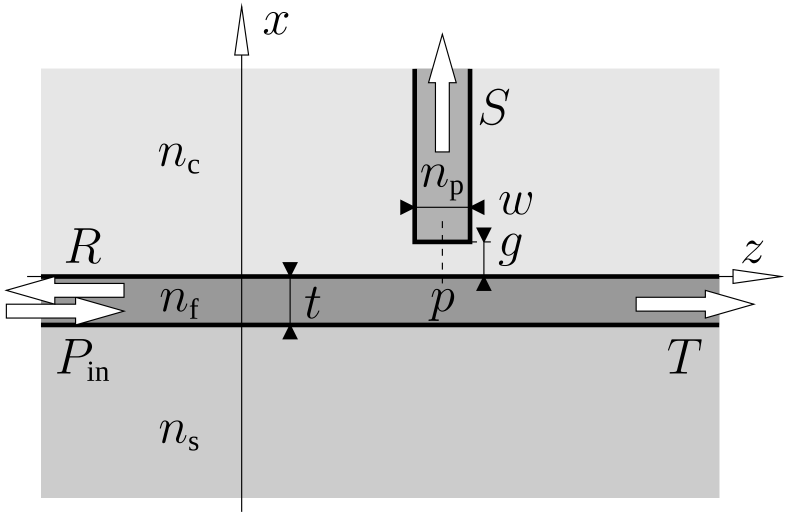
Probing evanescent fields

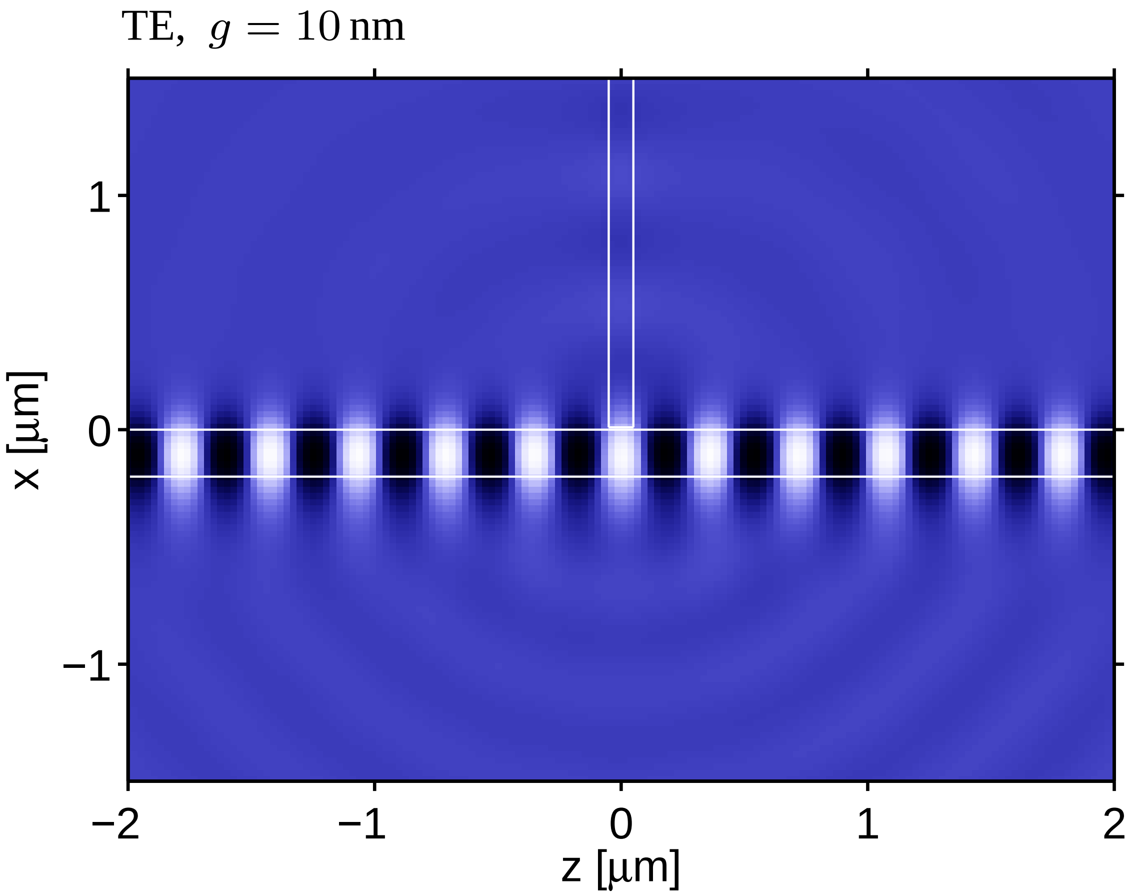




$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $w = 100\,$nm, $n_p = 1.5$,
$\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.0, 3.0]\,\mu\mathrm{m}^2$, $M_x = M_z = 80$.

Probing evanescent fields





$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $w = 100\,$nm, $n_p = 1.5$,
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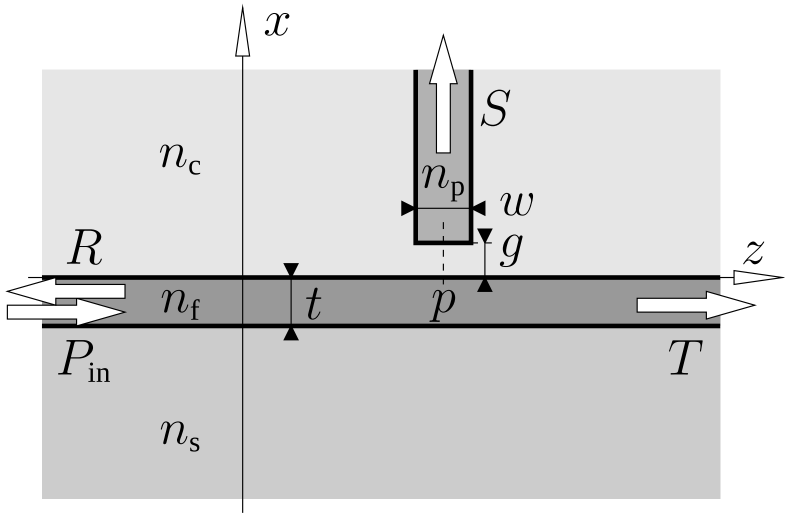
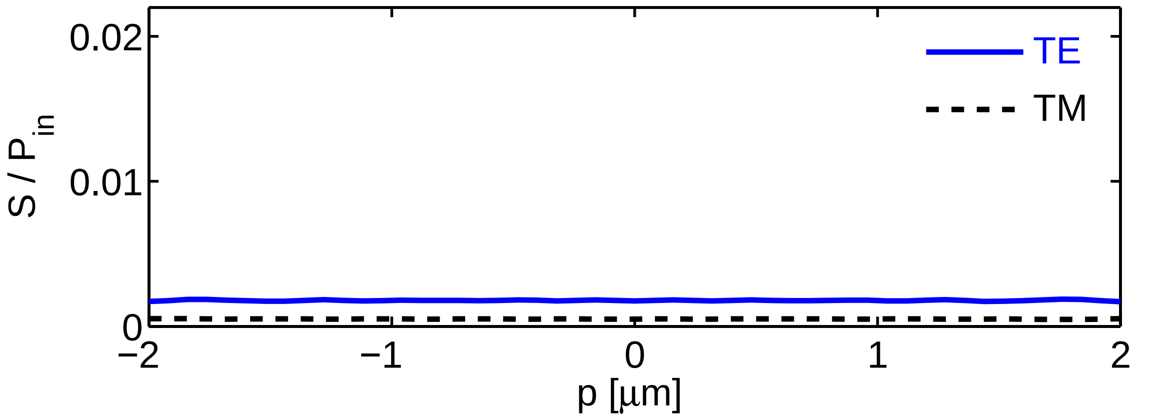
Probing evanescent fields



Horizontal scan: $g = 10\,\mathrm{nm}$



$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $w = 100\,$nm, $n_p = 1.5$,
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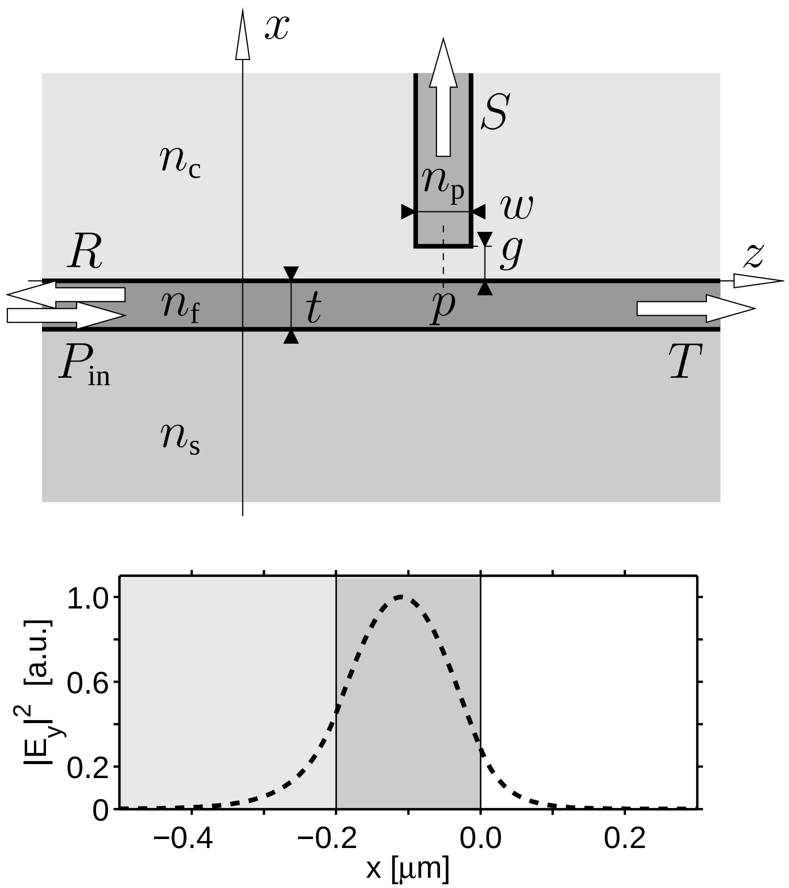
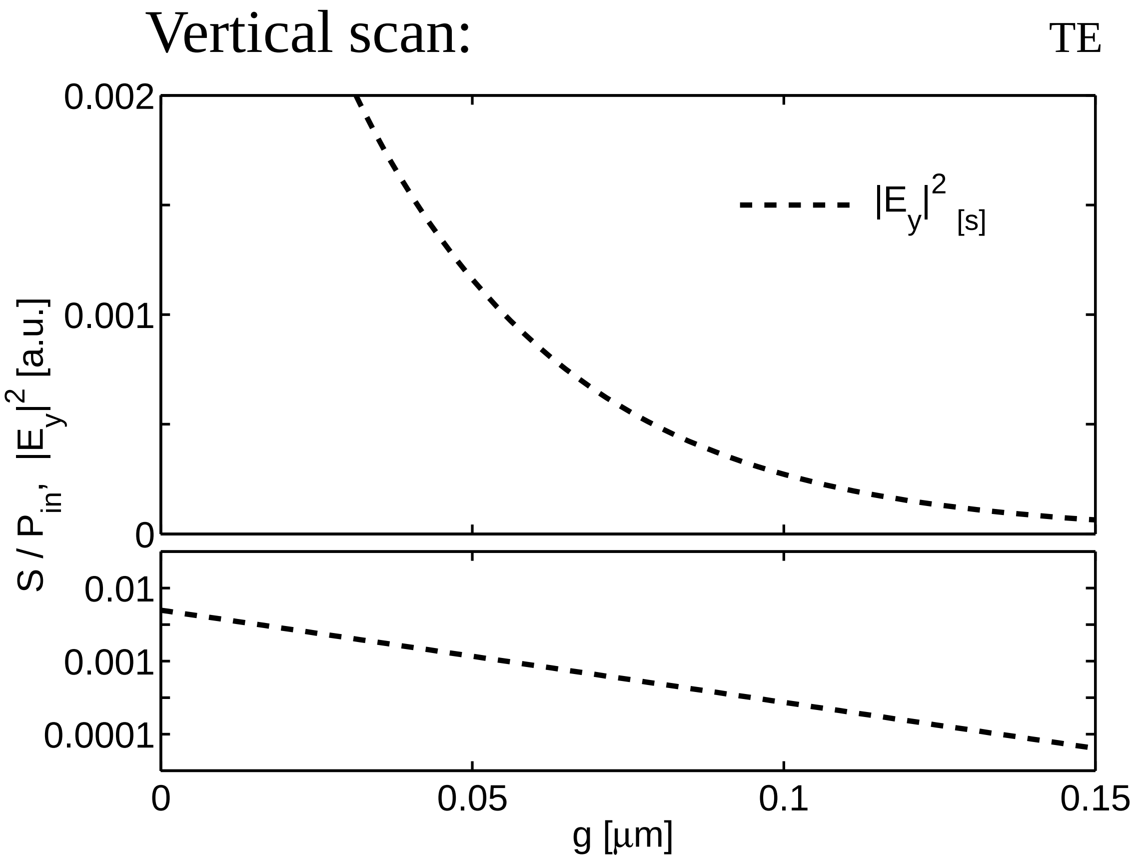
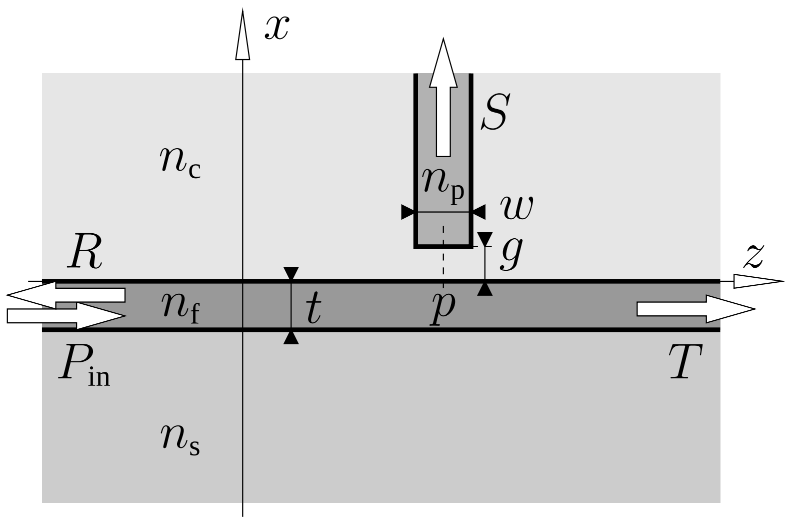
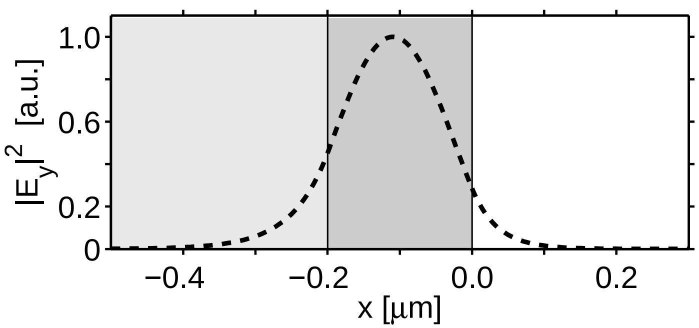
Probing evanescent fields



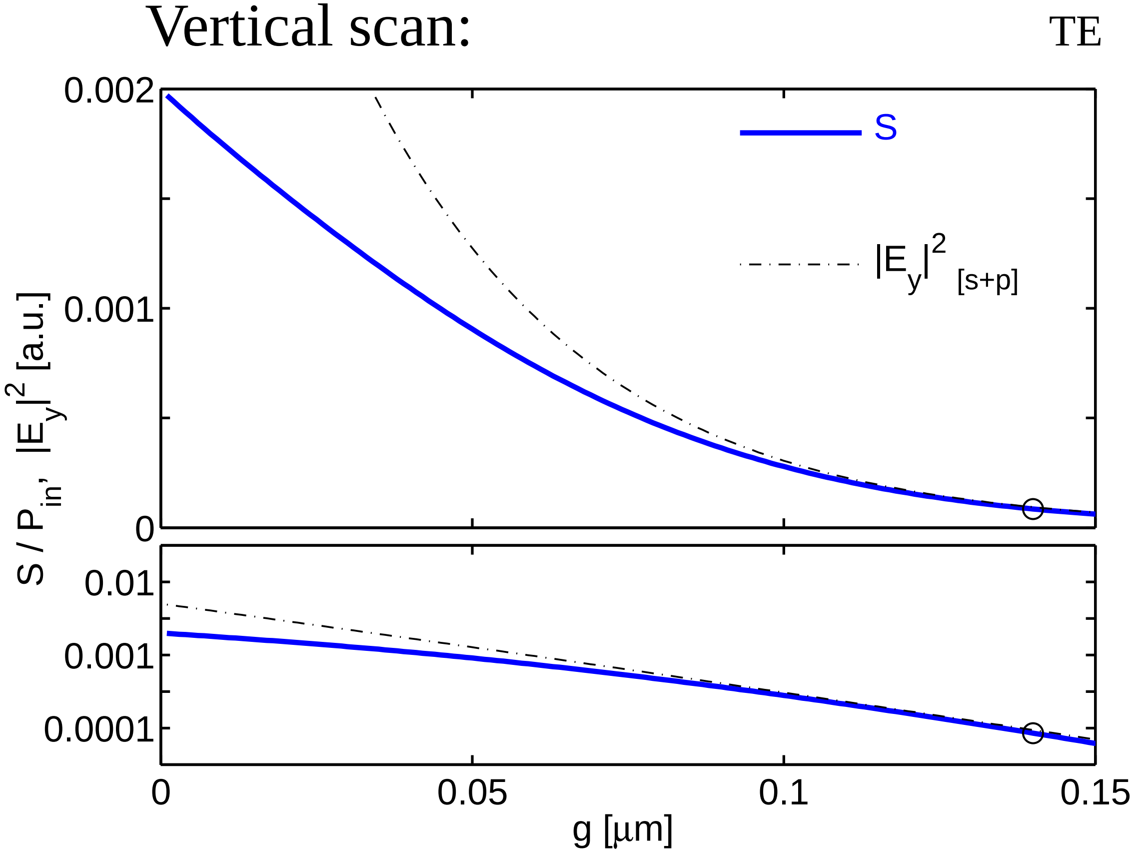


Probing evanescent fields

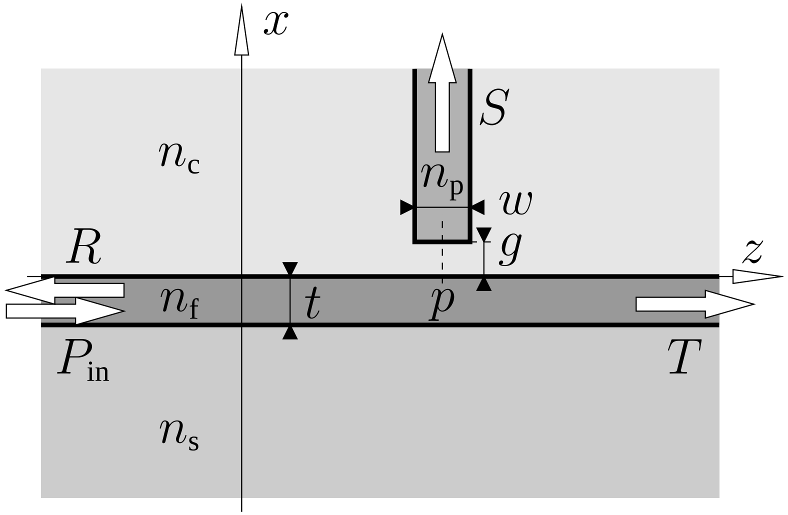


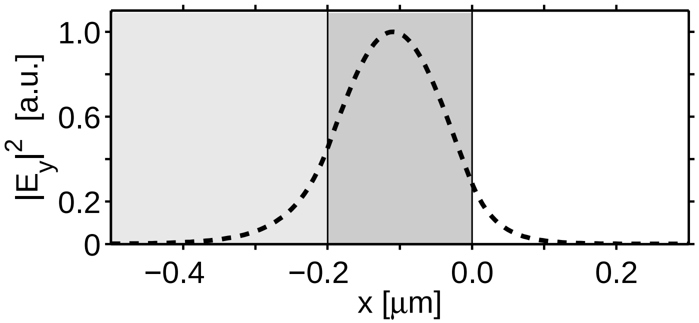


Vertical scan:



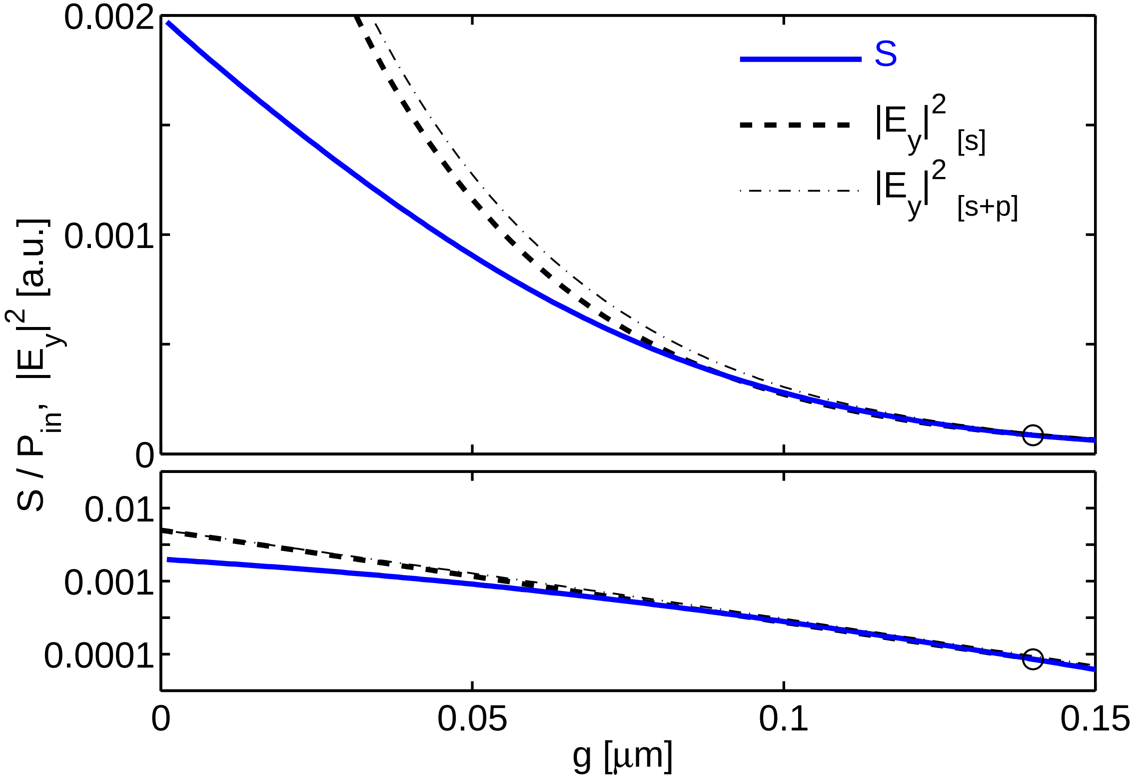
Probing evanescent fields



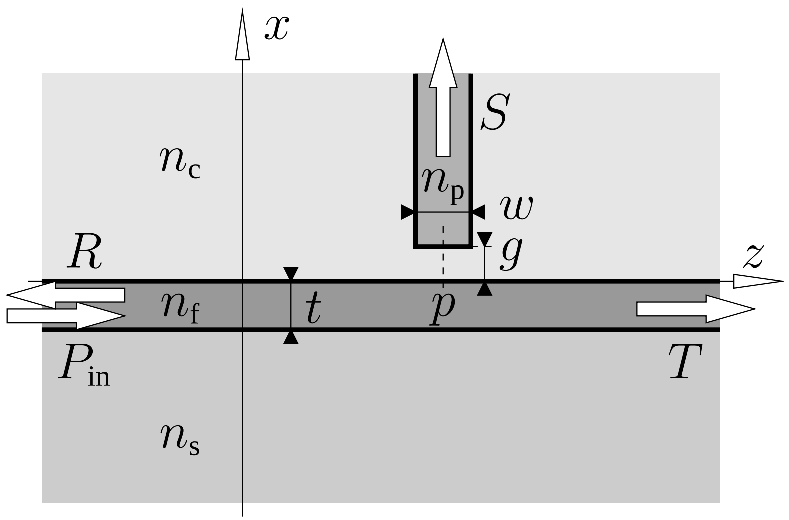


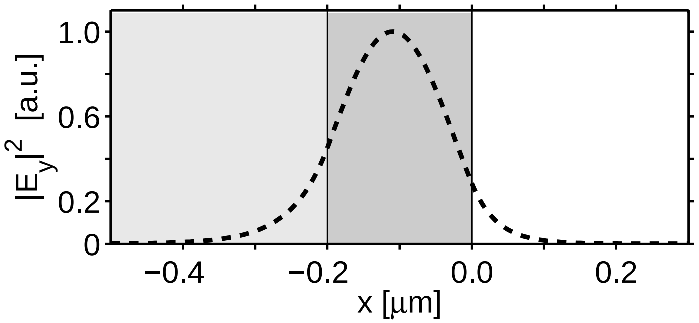
Vertical scan:

TE

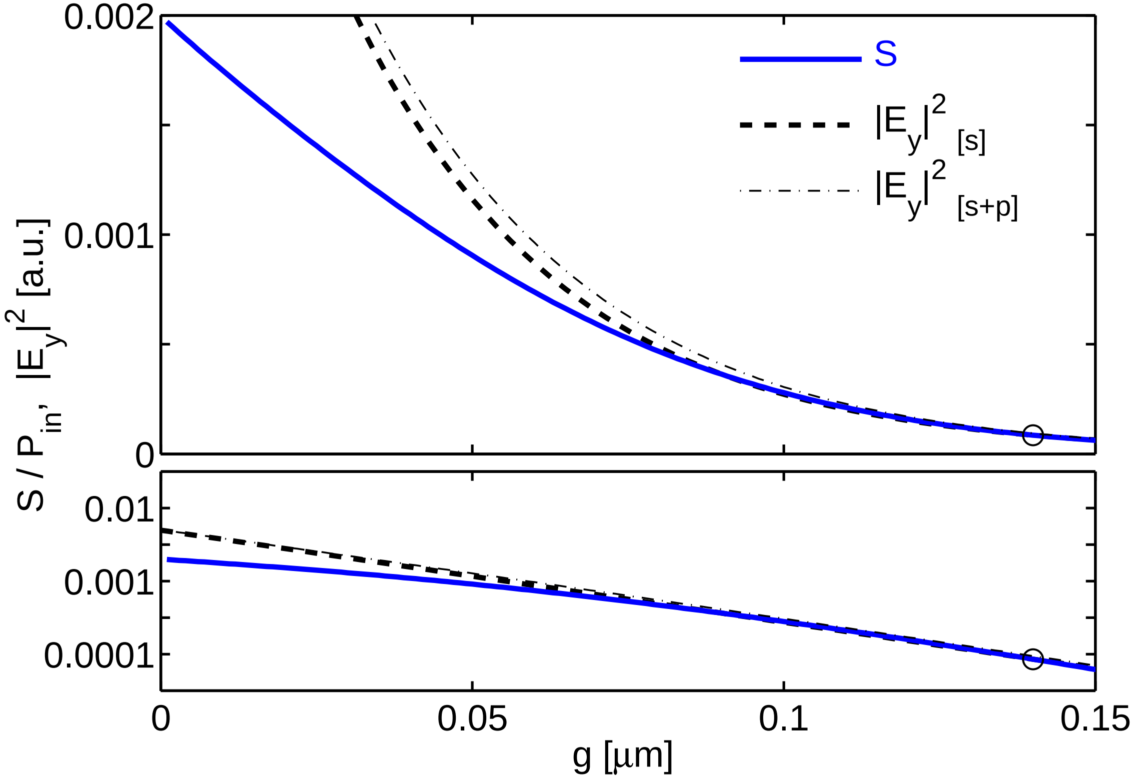


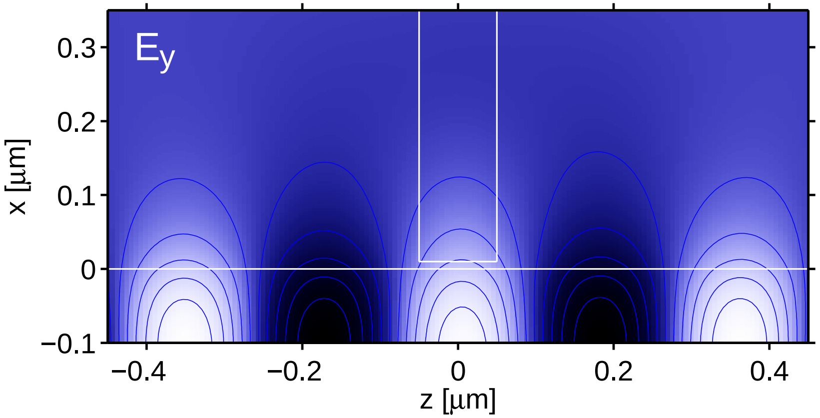
Probing evanescent fields

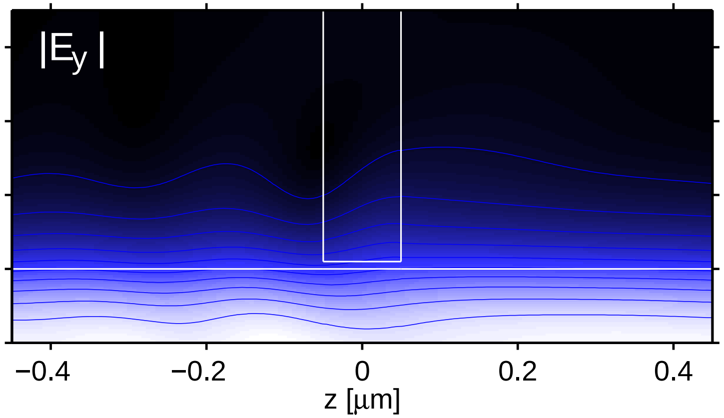




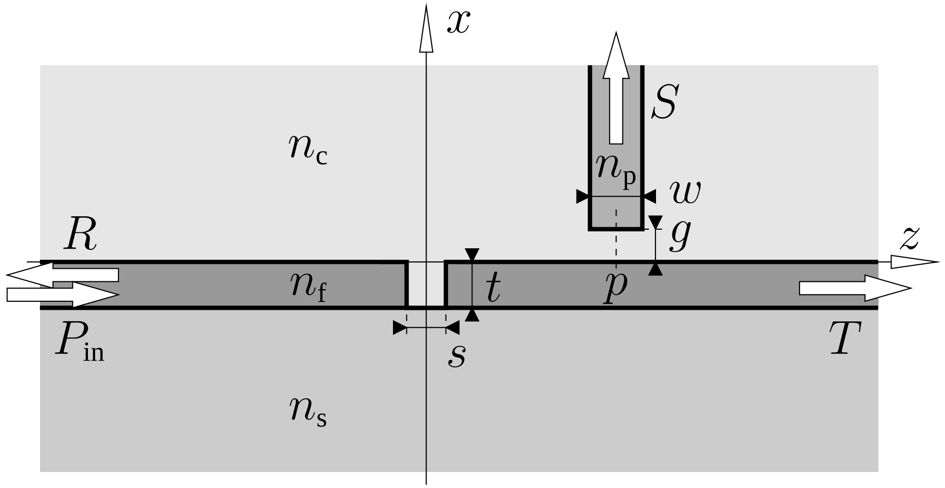
Vertical scan: TE

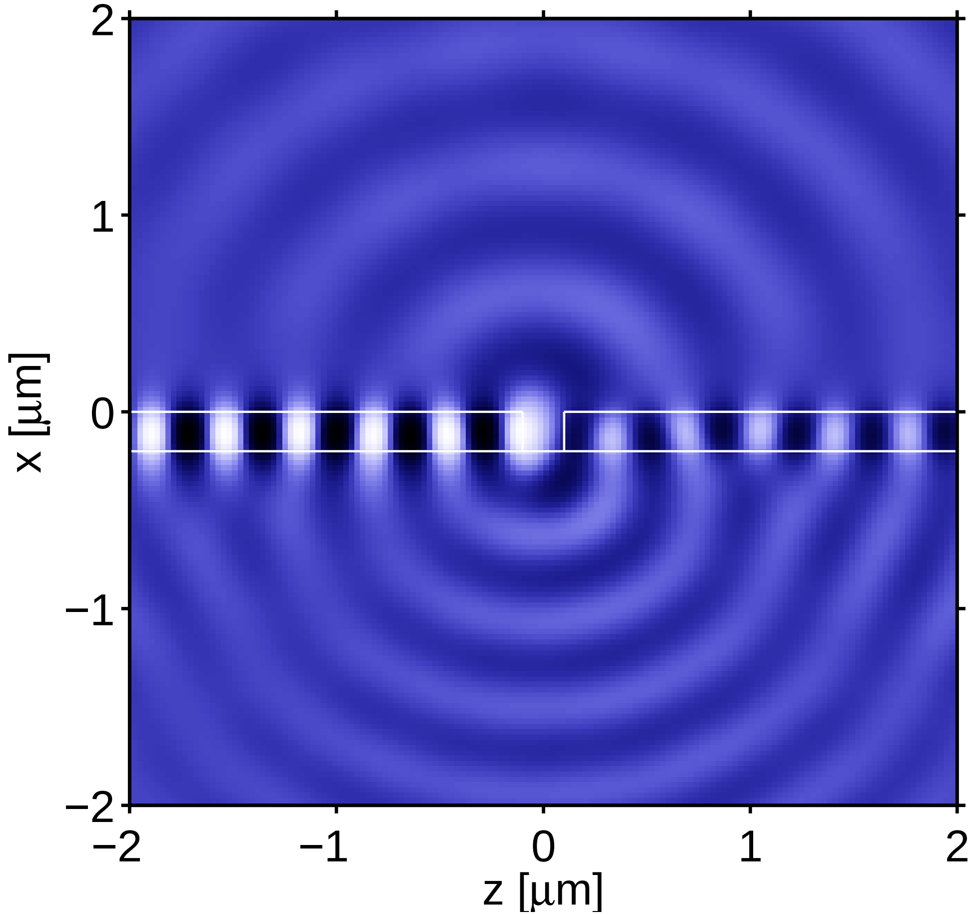






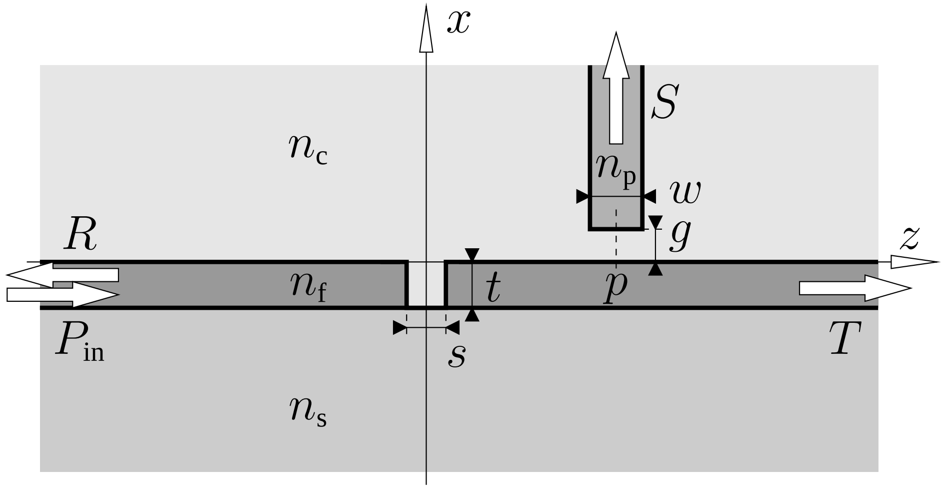
Hole defect in a slab waveguide

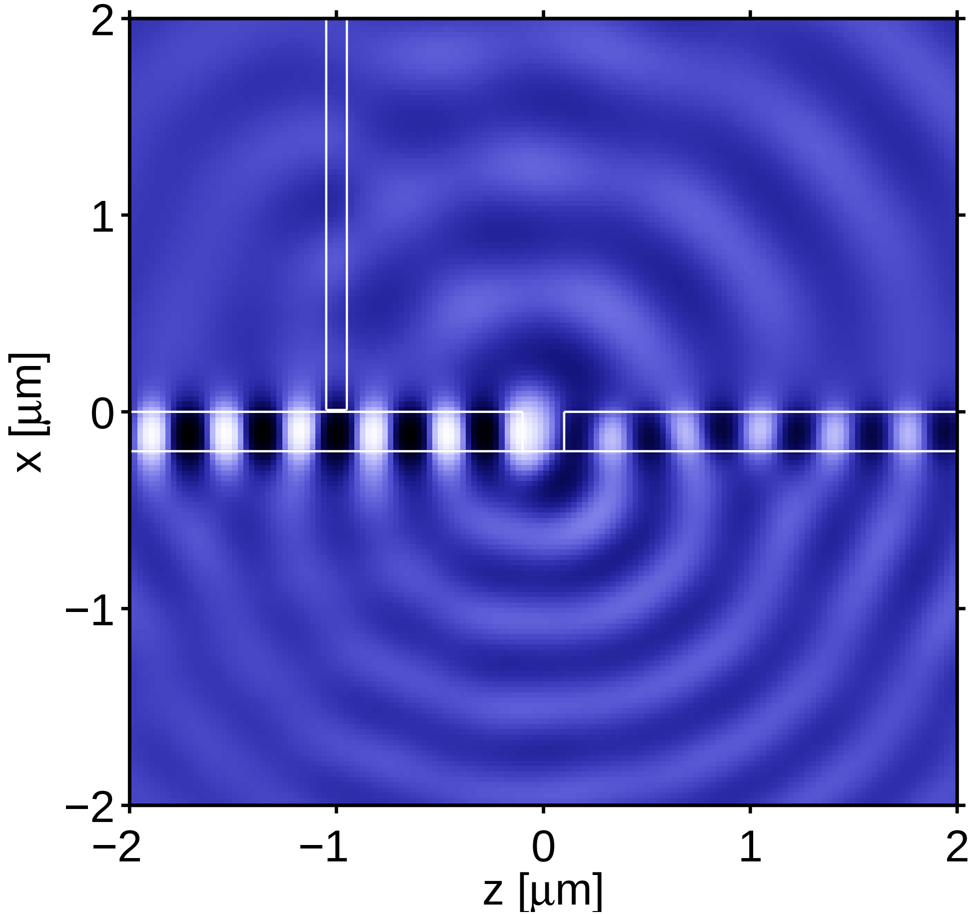




$n_\mathrm{s} = 1.45$, $n_\mathrm{f} = 2.0$, $n_\mathrm{c} = 1.0$, $t = 0.2\,\mu$m, $s = 0.2\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_\mathrm{p} = 1.5$,
TE, $\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.1, 3.1]\,\mu\mathrm{m}^2$, $M_x = M_z = 80$.

Hole defect in a slab waveguide





$n_{\rm s} = 1.45$, $n_{\rm f} = 2.0$, $n_{\rm c} = 1.0$, $t = 0.2\,\mu$m, $s = 0.2\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_{\rm p} = 1.5$,
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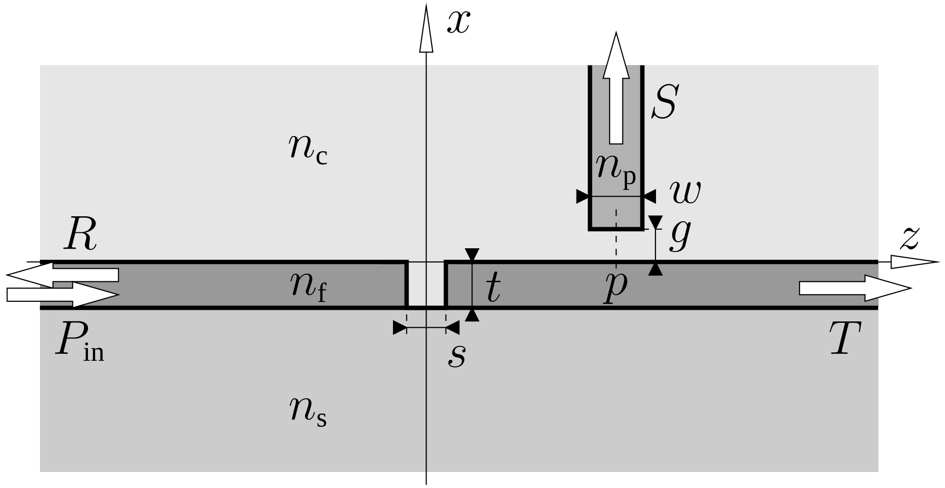
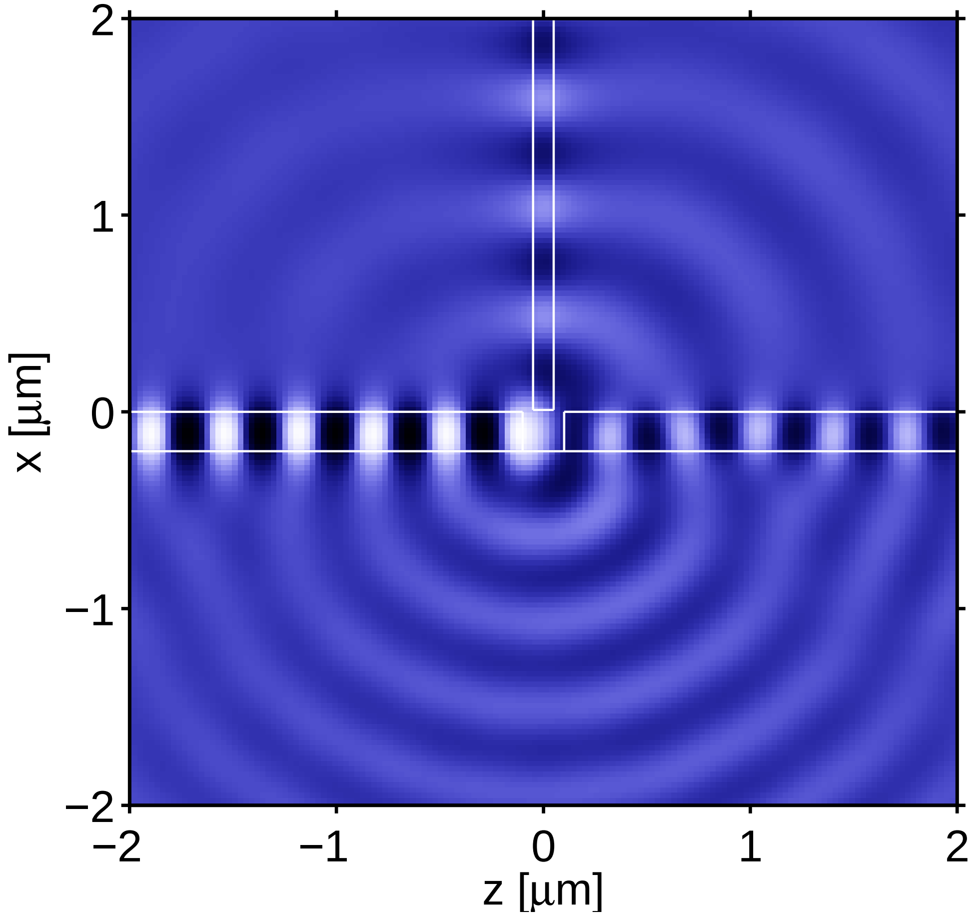
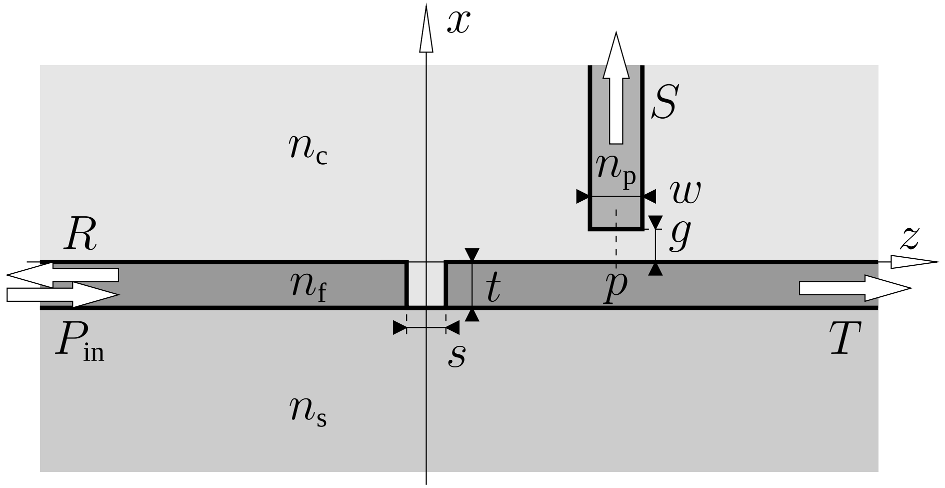
Hole defect in a slab waveguide

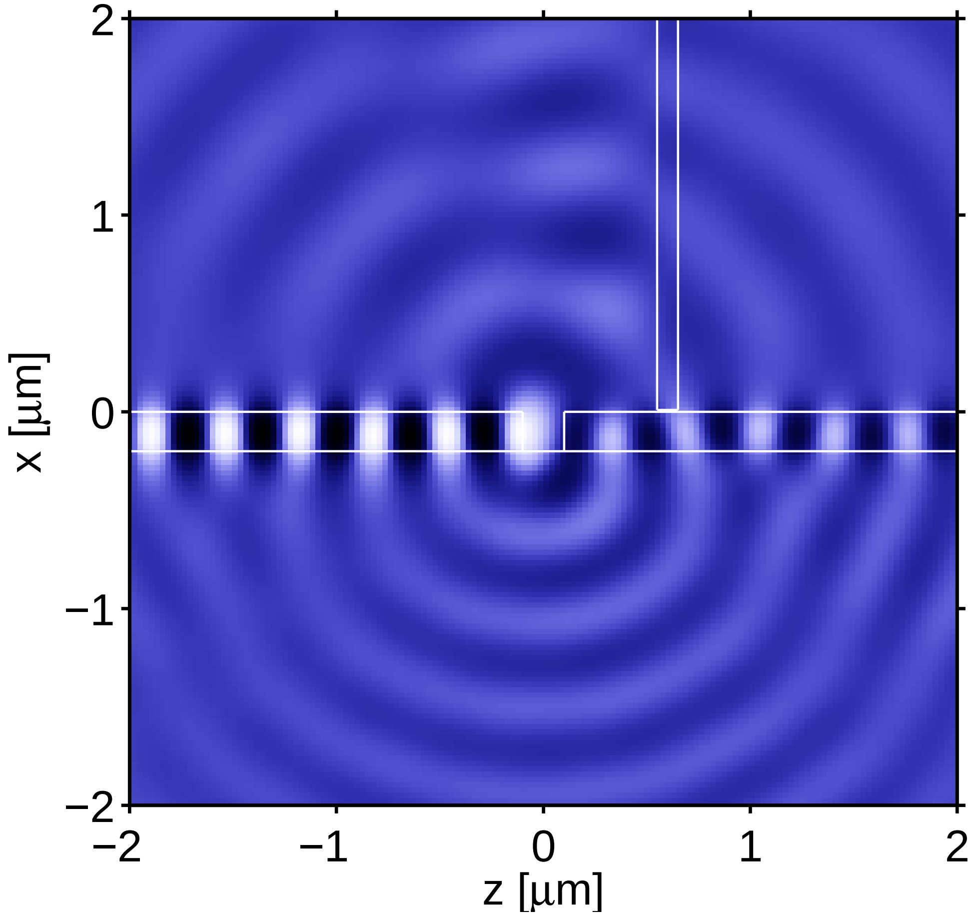




$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $s = 0.2\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_p = 1.5$,
TE, $\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.1, 3.1]\,\mu\text{m}^2$, $M_x = M_z = 80$.

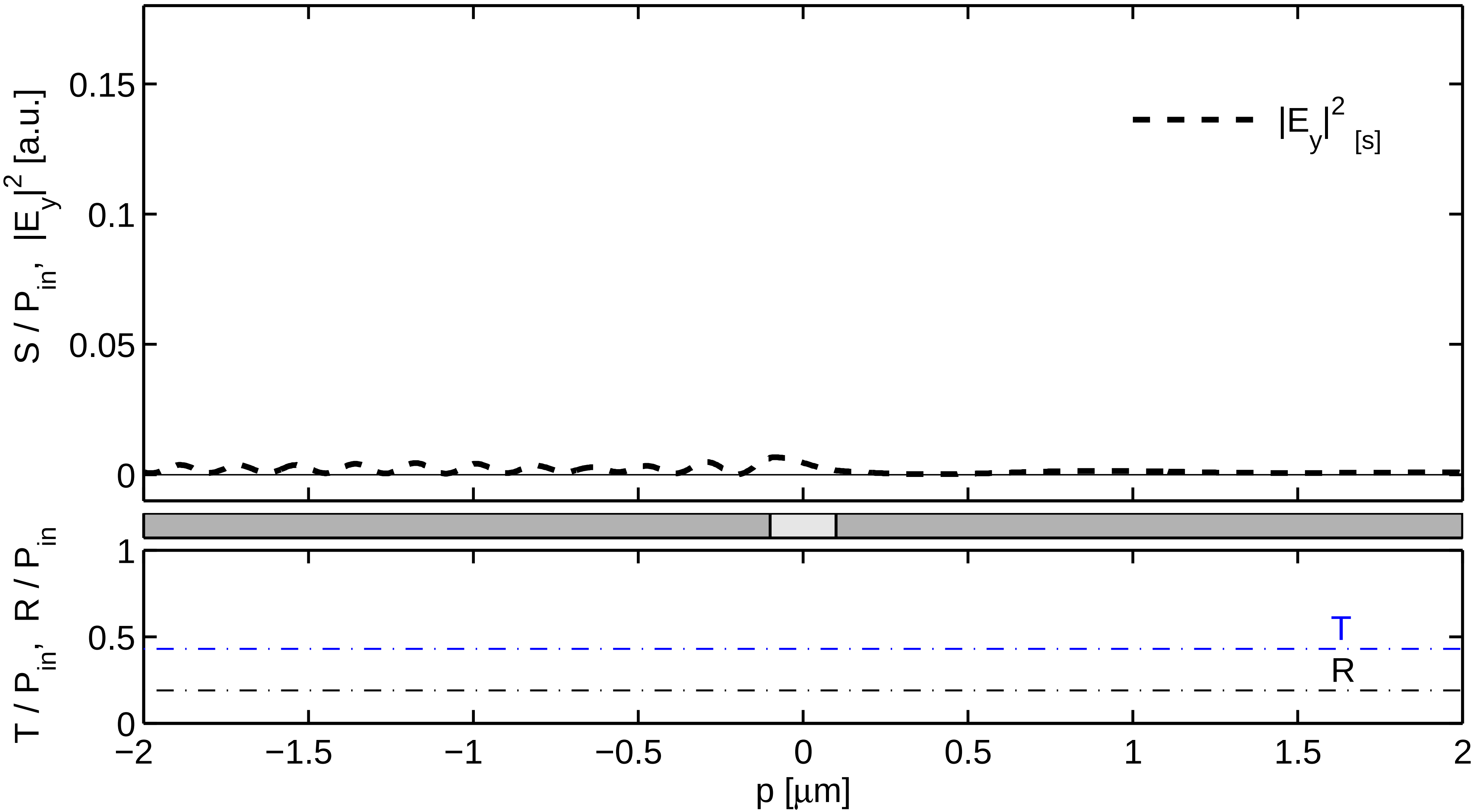
Hole defect in a slab waveguide



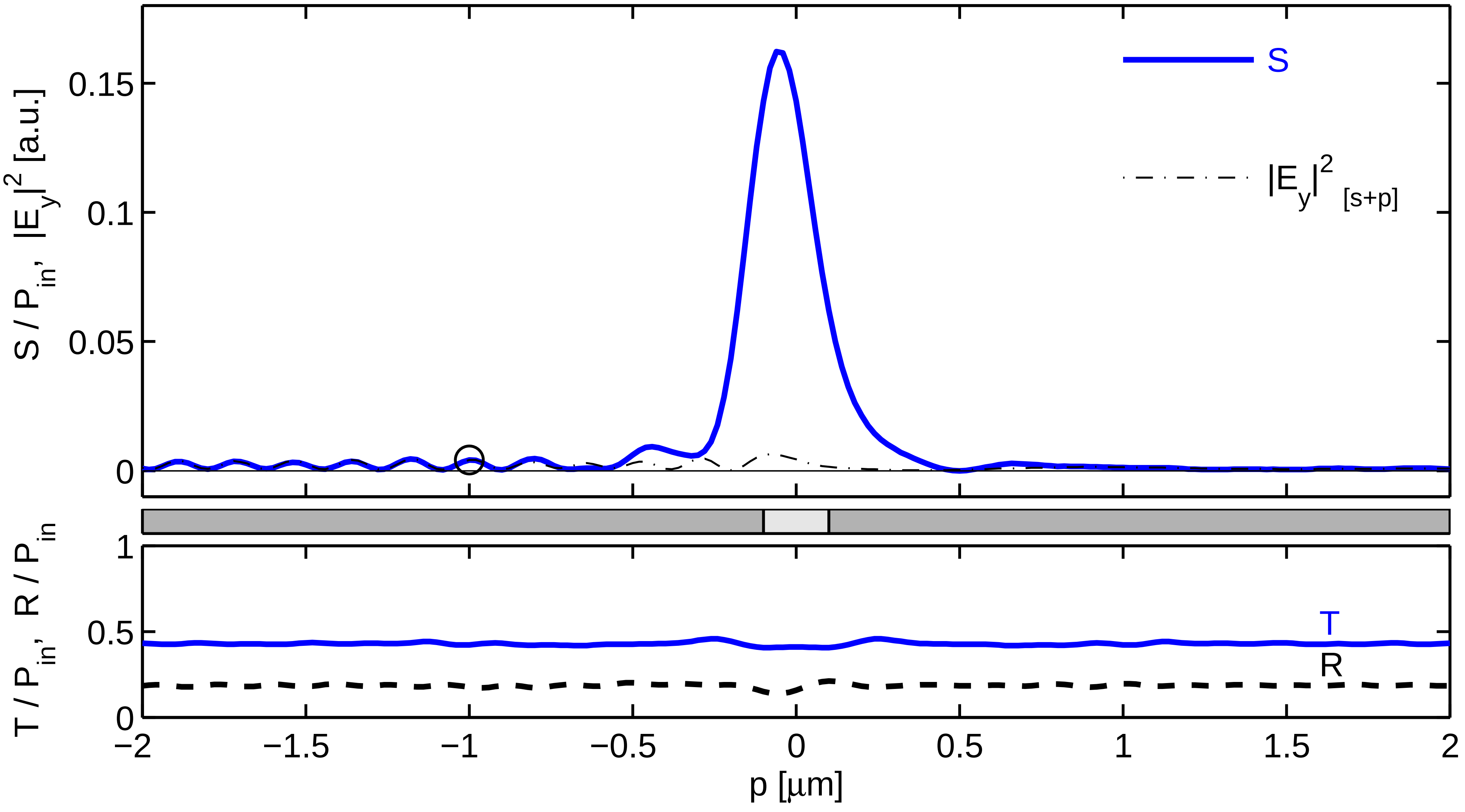


$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $s = 0.2\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_p = 1.5$,
TE, $\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.1, 3.1]\,\mu\text{m}^2$, $M_x = M_z = 80$.

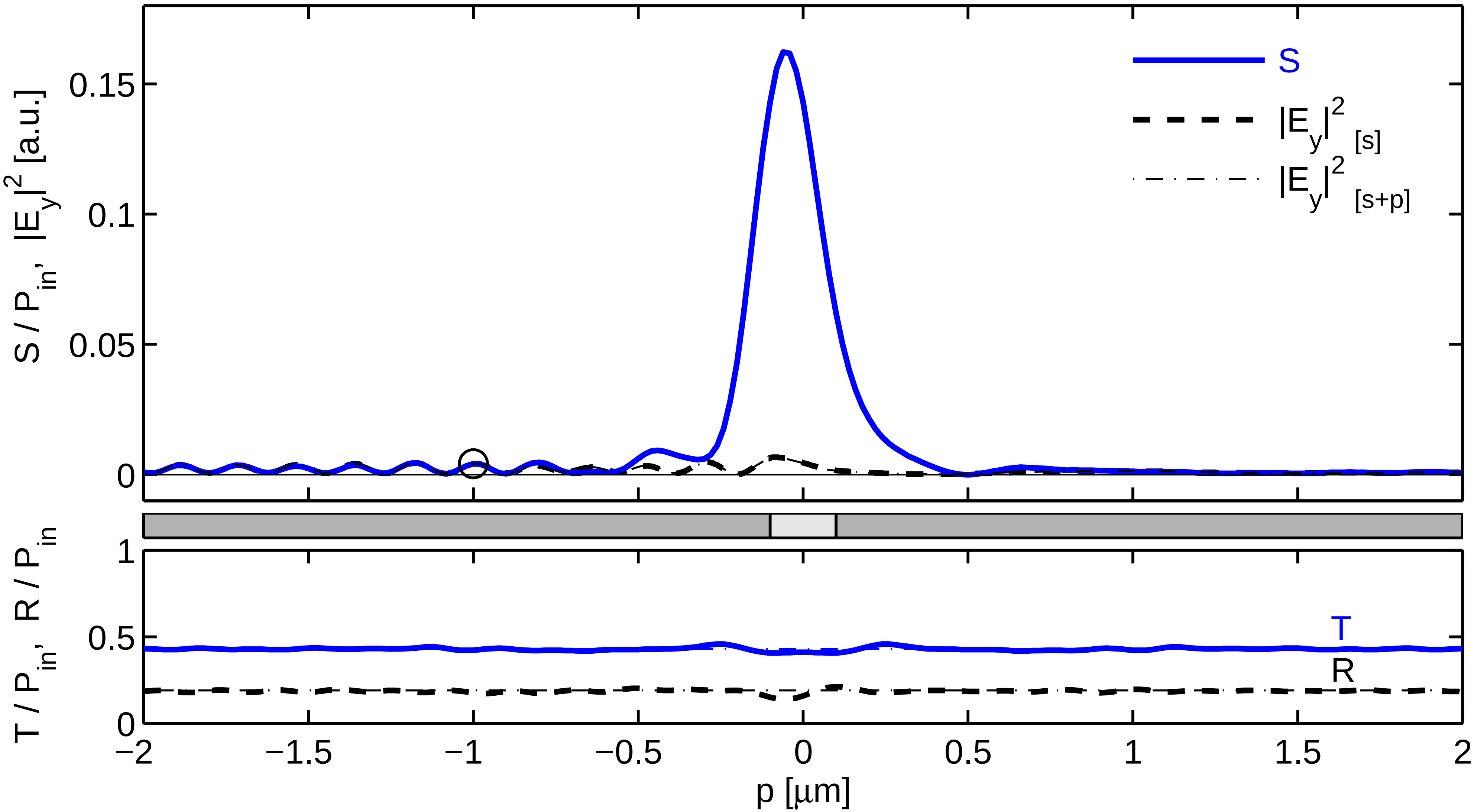
Hole defect in a slab waveguide



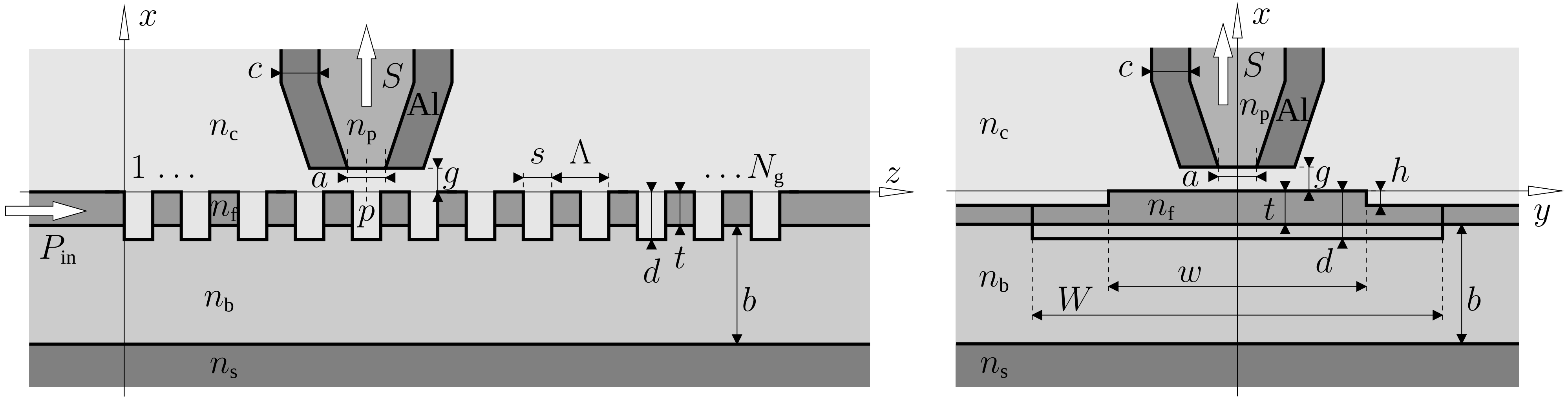
Hole defect in a slab waveguide



Hole defect in a slab waveguide



Bragg grating: PSTM experiment *

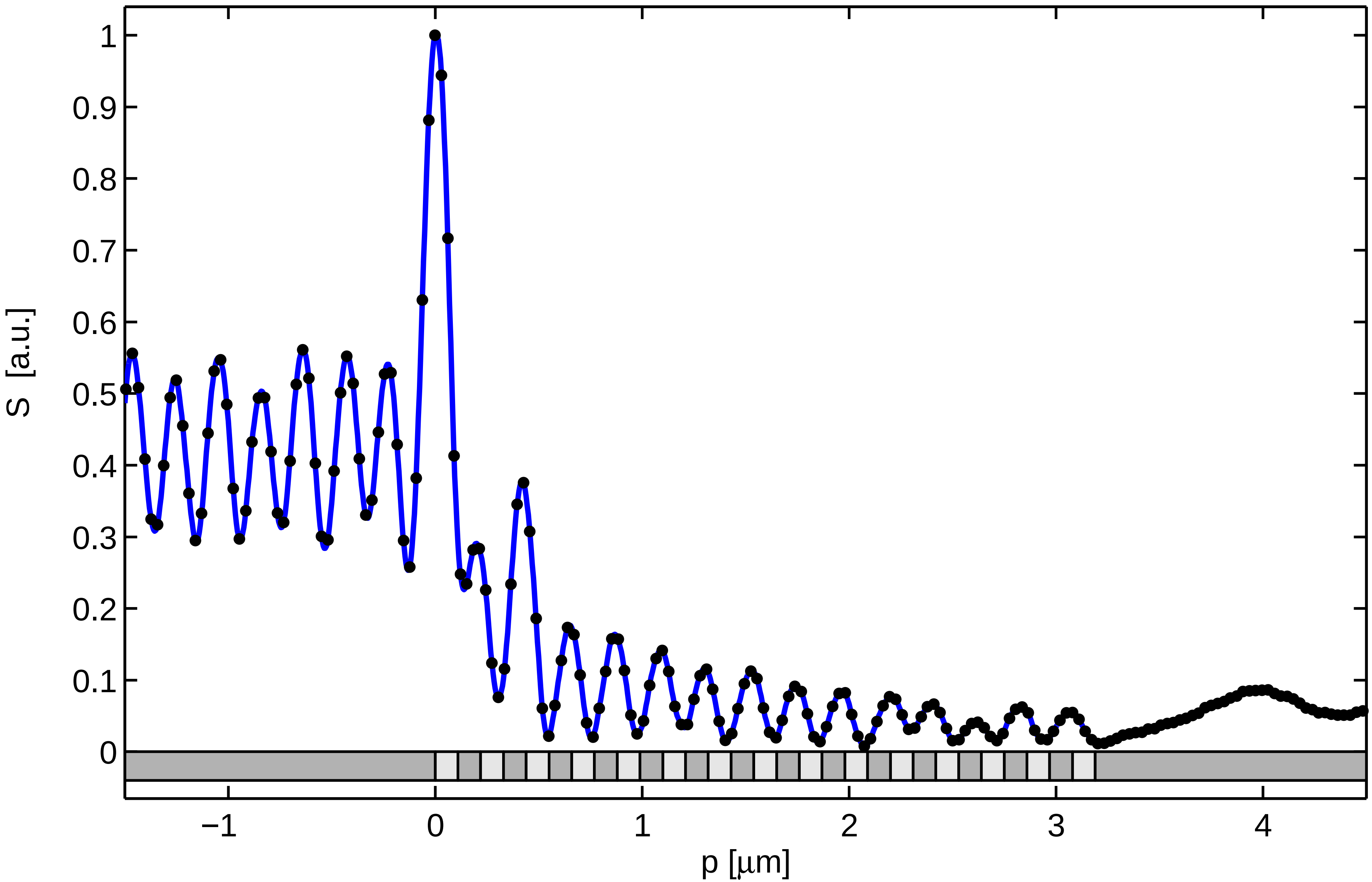


Sample: Rib waveguide with a series of deep, rectangular slits,
$n_s = 3.4$, $n_b = 1.45$, $n_f = 2.01$, $n_c = 1.0$, $t = 55\,$nm, $h = 11\,$nm, $w = 1.5\,\mu$m,
$W = 2.5\,\mu$m, $b = 3.2\,\mu$m, $\Lambda = 220\,$nm, $s = 110\,$nm, $d = 70\,$nm, $N_g = 15$.

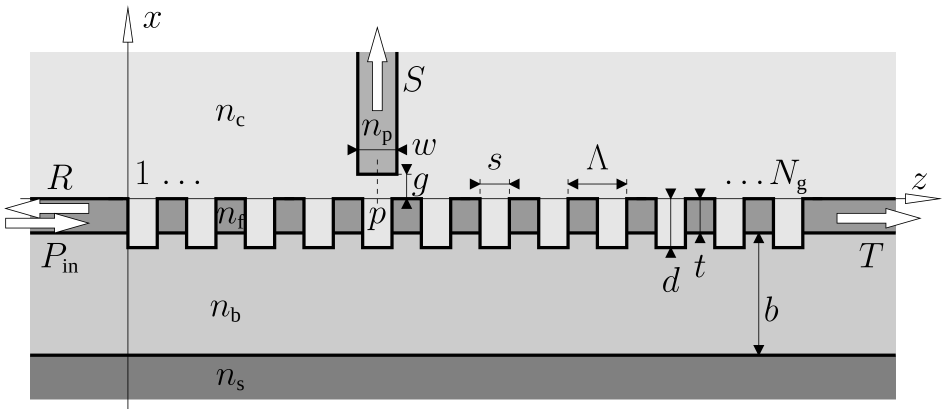
Probe: Tapered cylindrical fiber tip with aluminium coating,
$a \approx 80\,$nm, $c \approx 100\,$nm, $g = 10\,$nm, $n_p = 1.5$;
TE polarized light, vacuum wavelength $\lambda = 0.6328\,\mu$m.

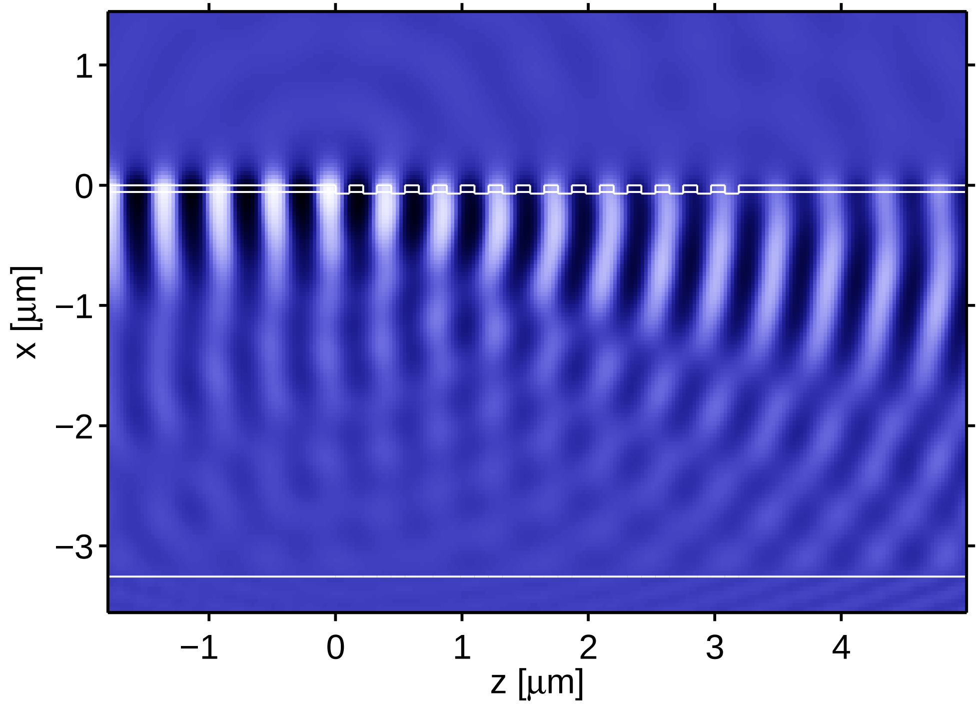
* E. Flück, M. Hammer, A. M. Otter, J. P. Korterijk, L. Kuipers, N. F. van Hulst,
Amplitude and phase evolution of optical fields inside periodic photonic structures,
Journal of Lightwave Technology **21**(5), 1384-1393 (2003)

Bragg grating: PSTM experiment



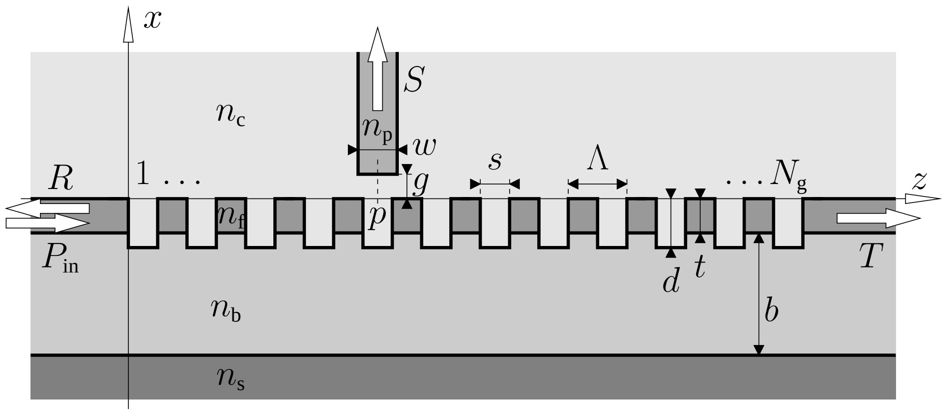
Waveguide Bragg grating

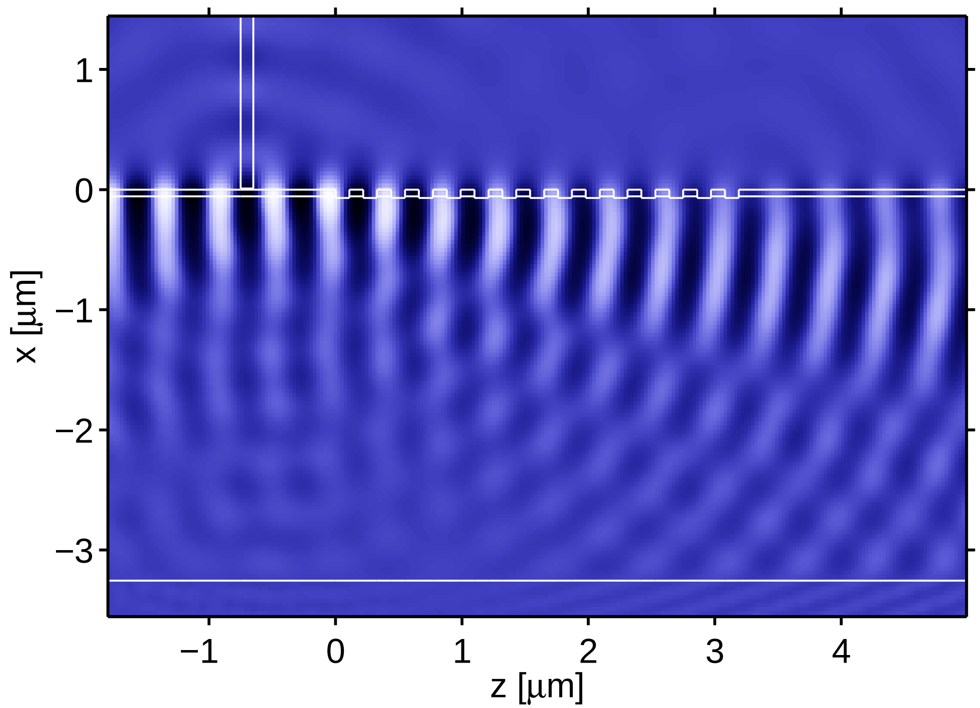




$n_s = 3.4$, $n_b = 1.45$, $n_f = 2.01$, $n_c = 1.0$, $t = 55\,\mathrm{nm}$, $b = 3.2\,\mu\mathrm{m}$, $d = 70\,\mathrm{nm}$,
$\Lambda = 220\,\mathrm{nm}$, $s = 110\,\mathrm{nm}$, $N_g = 15$, $g = 10\,\mathrm{nm}$, $w = 100\,\mathrm{nm}$, $n_p = 1.5$,
TE, $\lambda = 0.6328\,\mu\mathrm{m}$, $(x, z) \in [-3.5, 1.5] \times [-2.0, 5.2]\,\mu\mathrm{m}^2$, $M_x = 80$, $M_z = 100$.

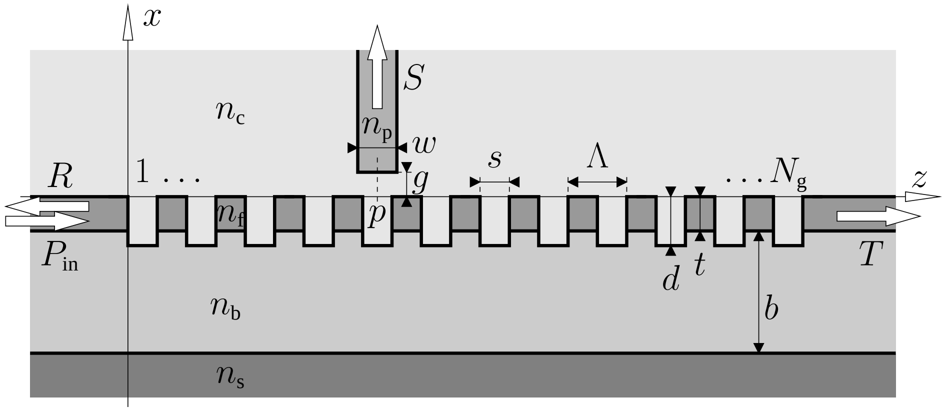
Waveguide Bragg grating

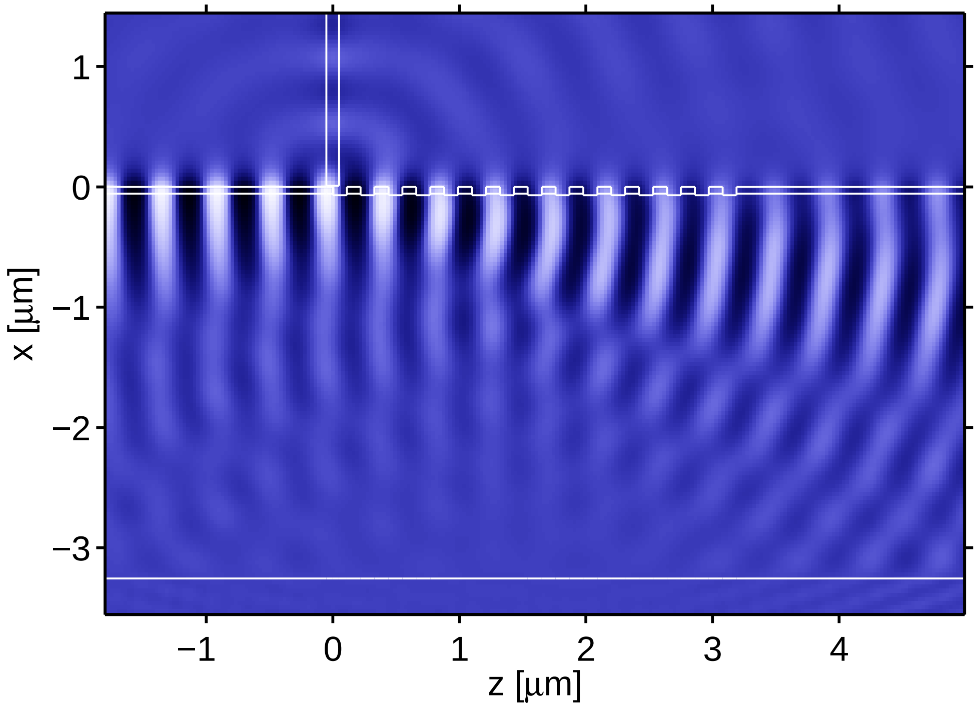




$n_s = 3.4$, $n_b = 1.45$, $n_f = 2.01$, $n_c = 1.0$, $t = 55\,\text{nm}$, $b = 3.2\,\mu\text{m}$, $d = 70\,\text{nm}$,
$\Lambda = 220\,\text{nm}$, $s = 110\,\text{nm}$, $N_g = 15$, $g = 10\,\text{nm}$, $w = 100\,\text{nm}$, $n_p = 1.5$,
TE, $\lambda = 0.6328\,\mu\text{m}$, $(x, z) \in [-3.5, 1.5] \times [-2.0, 5.2]\,\mu\text{m}^2$, $M_x = 80$, $M_z = 100$.

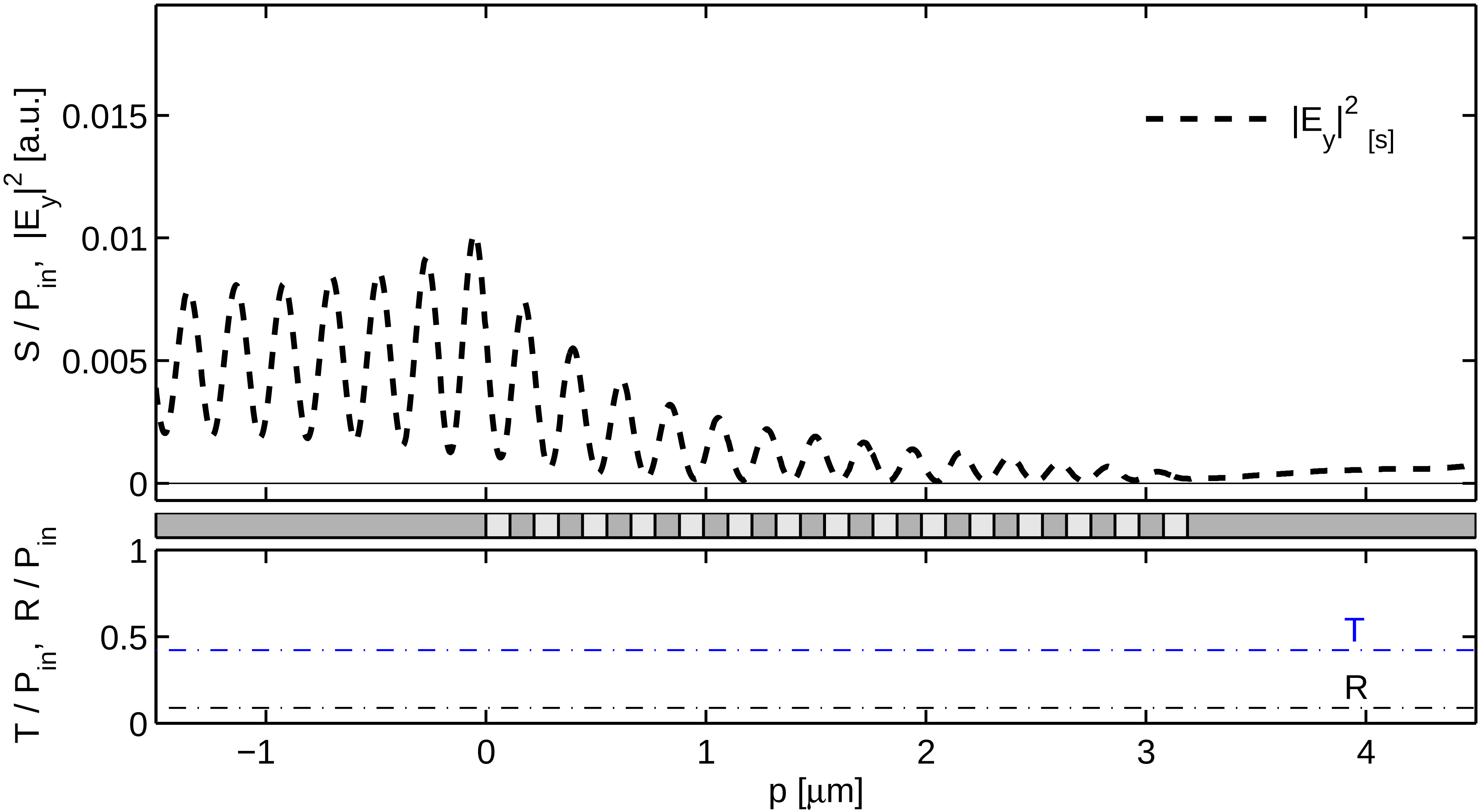
Waveguide Bragg grating



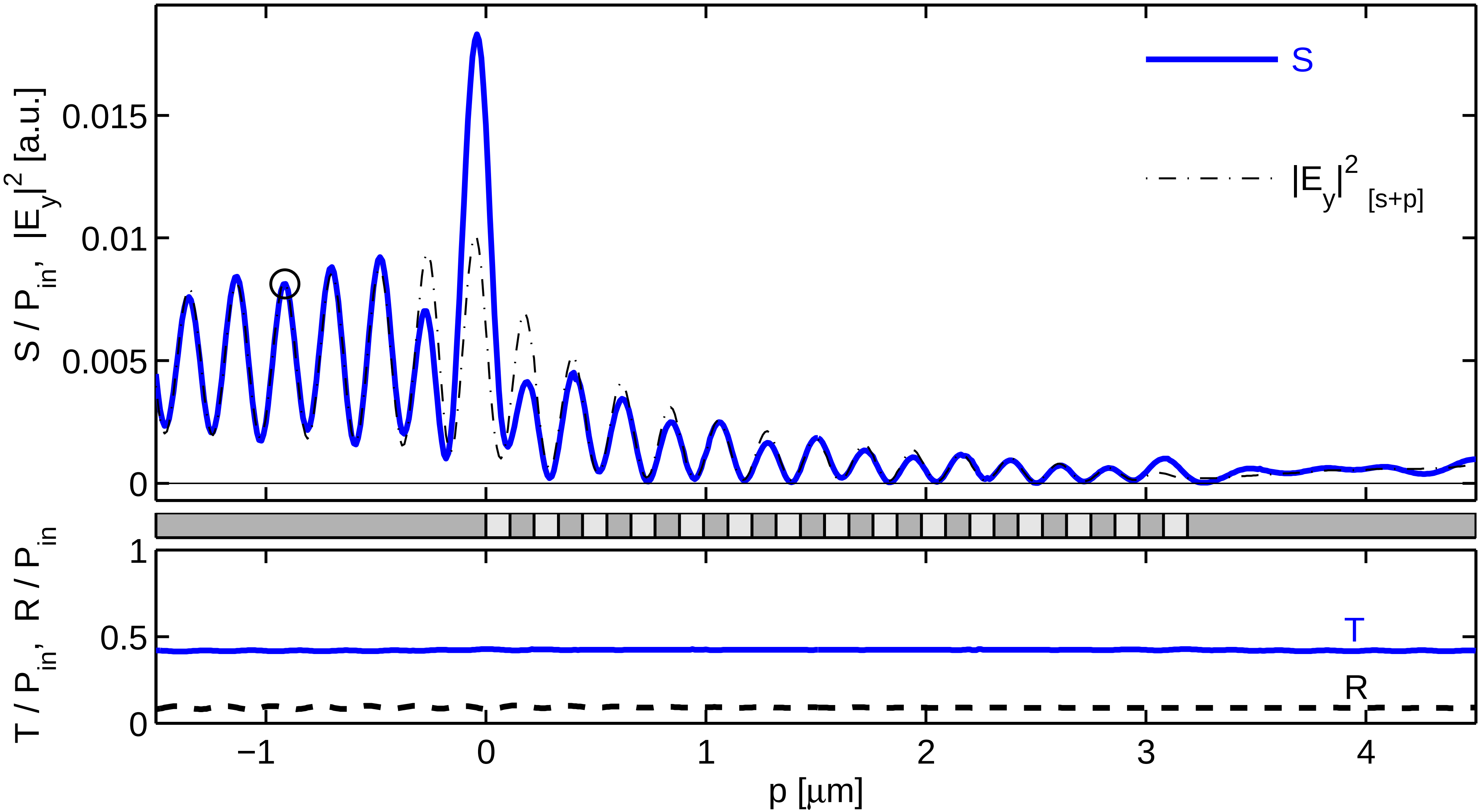


$n_s = 3.4$, $n_b = 1.45$, $n_f = 2.01$, $n_c = 1.0$, $t = 55\,\text{nm}$, $b = 3.2\,\mu\text{m}$, $d = 70\,\text{nm}$,
$\Lambda = 220\,\text{nm}$, $s = 110\,\text{nm}$, $N_g = 15$, $g = 10\,\text{nm}$, $w = 100\,\text{nm}$, $n_p = 1.5$,
TE, $\lambda = 0.6328\,\mu\text{m}$, $(x, z) \in [-3.5, 1.5] \times [-2.0, 5.2]\,\mu\text{m}^2$, $M_x = 80$, $M_z = 100$.

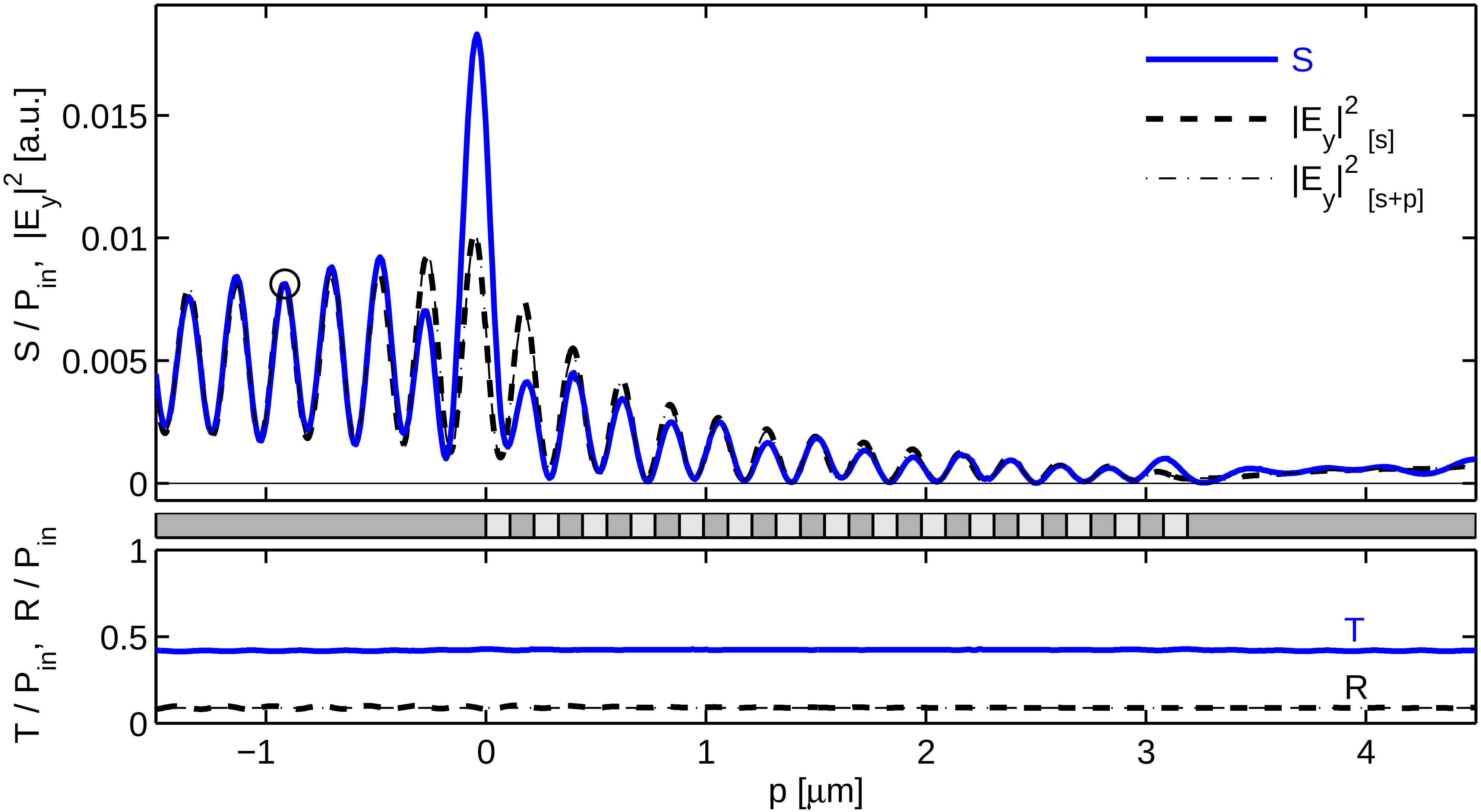
Waveguide Bragg grating



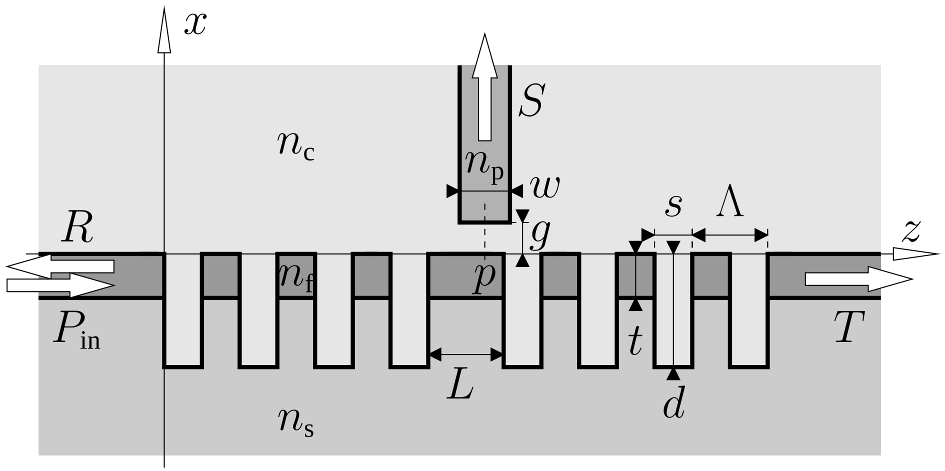
Waveguide Bragg grating

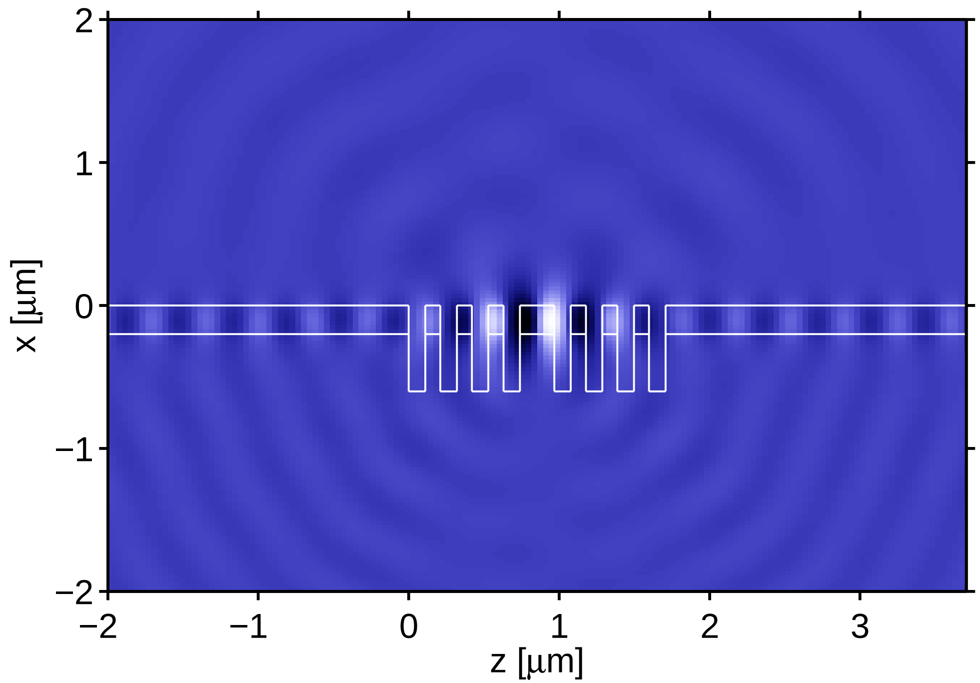


Waveguide Bragg grating



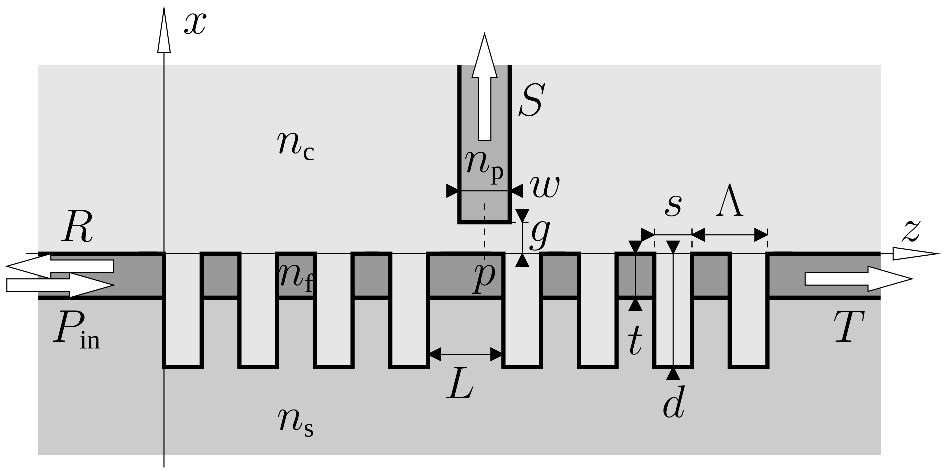
Resonant defect cavity

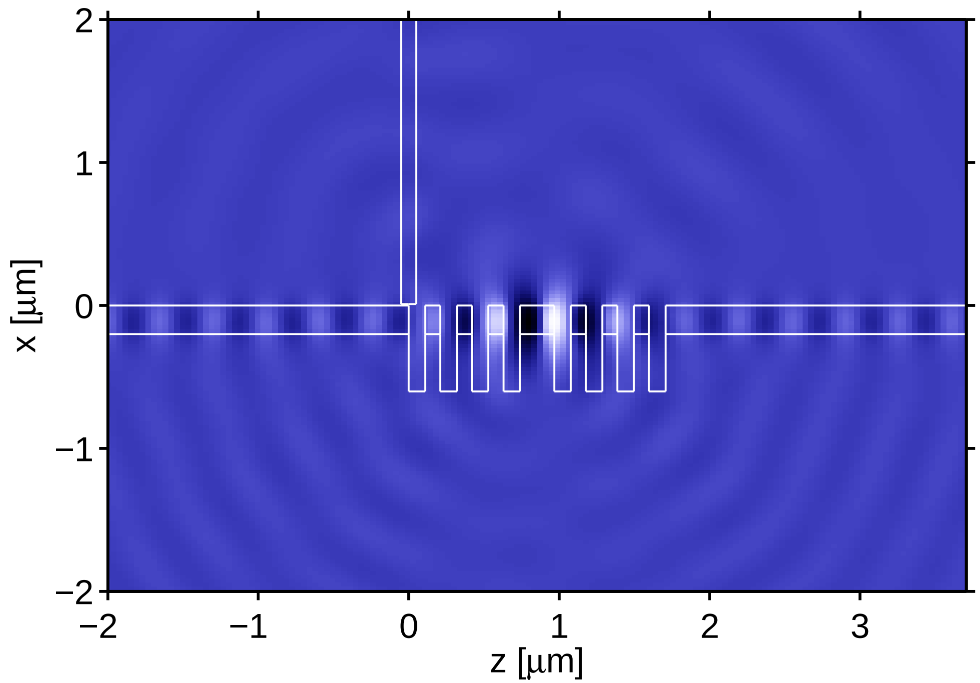




$n_\mathrm{s} = 1.45$, $n_\mathrm{f} = 2.0$, $n_\mathrm{c} = 1.0$, $t = 0.2\,\mu$m, $d = 0.6\,\mu$m,
$\Lambda = 0.21\,\mu$m, $s = 0.11\,\mu$m, $L = 0.2275\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_\mathrm{p} = 1.5$,
TE, $\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.0, 4.7075]\,\mu\mathrm{m}^2$, $M_x = 100$, $M_z = 120$.

Resonant defect cavity





$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $d = 0.6\,\mu$m,
$\Lambda = 0.21\,\mu$m, $s = 0.11\,\mu$m, $L = 0.2275\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_p = 1.5$,
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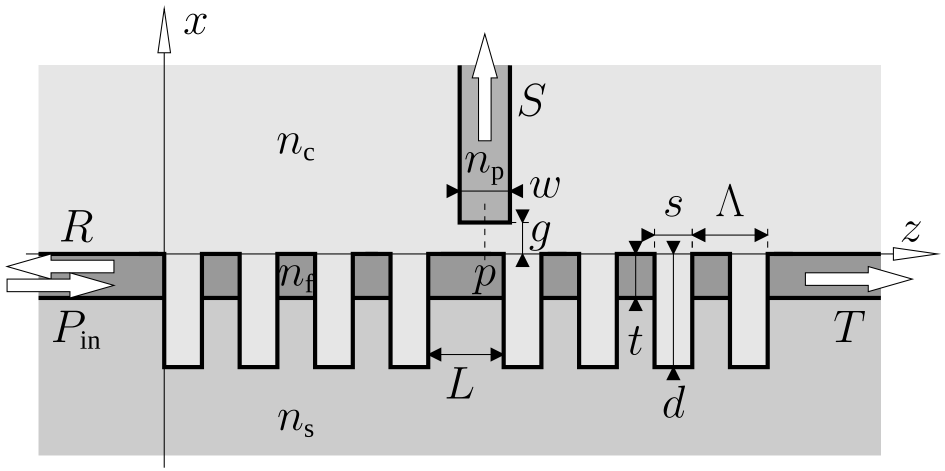
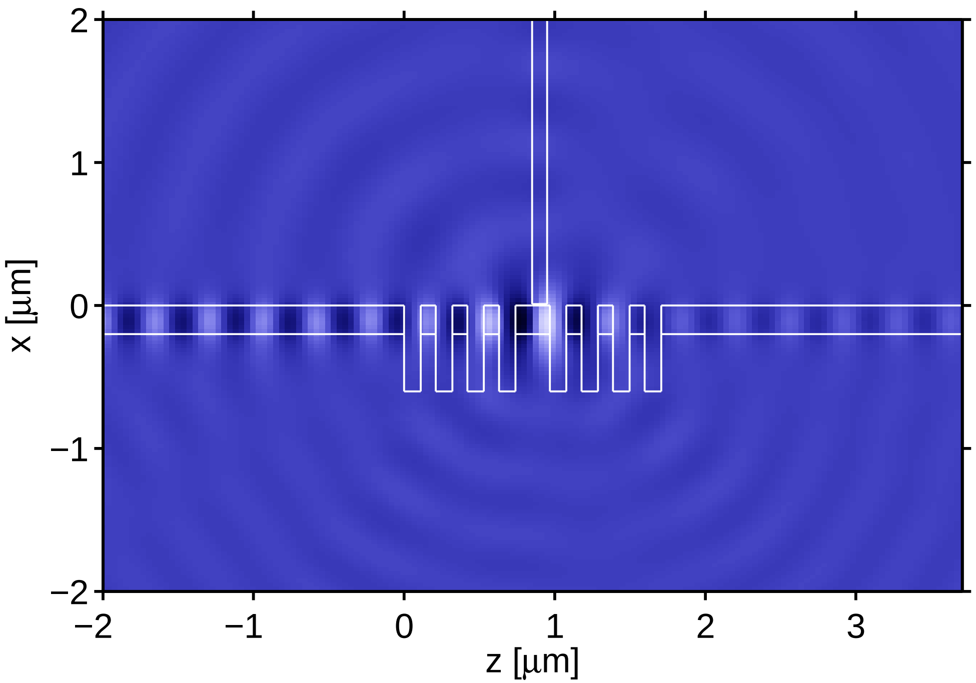
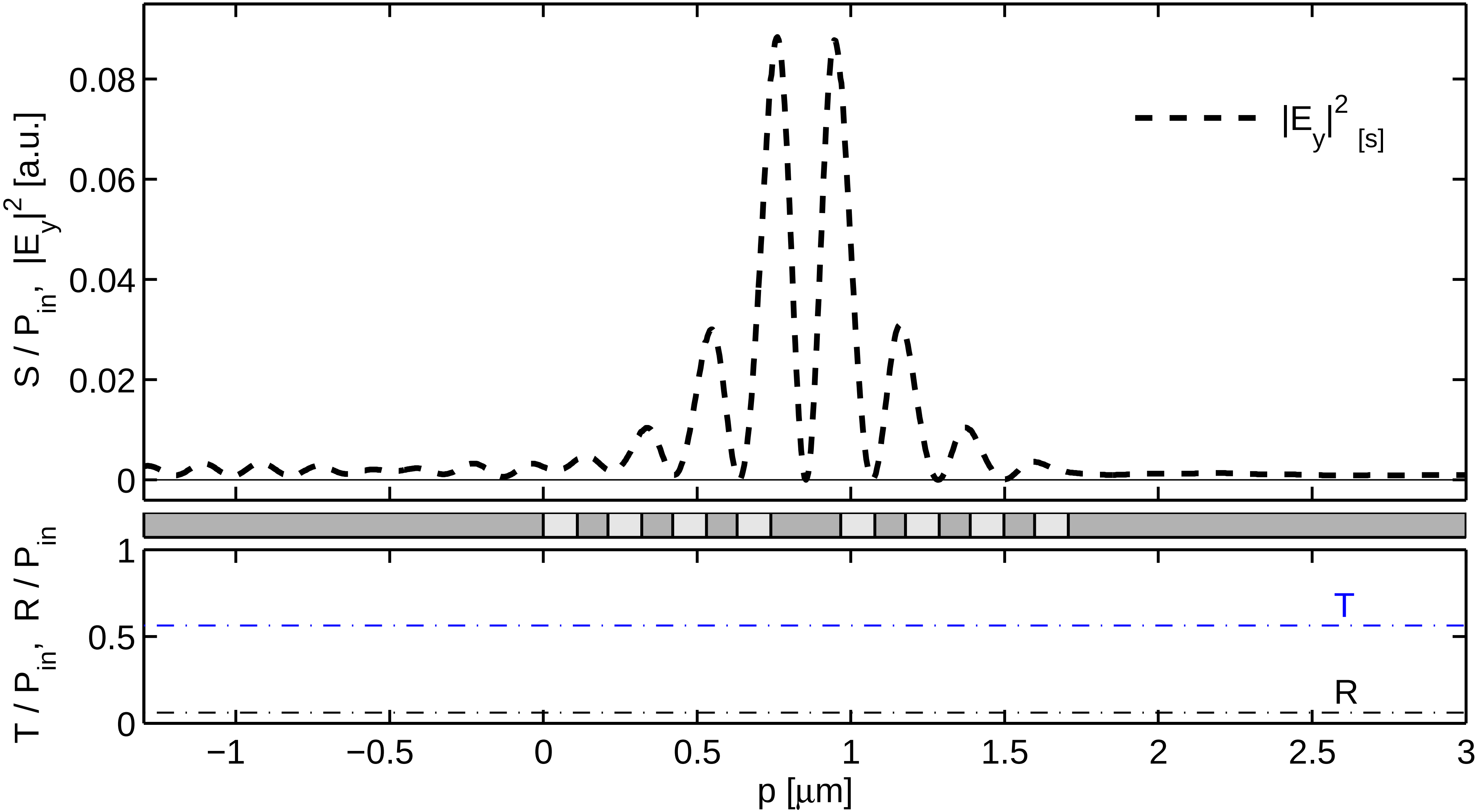
Resonant defect cavity



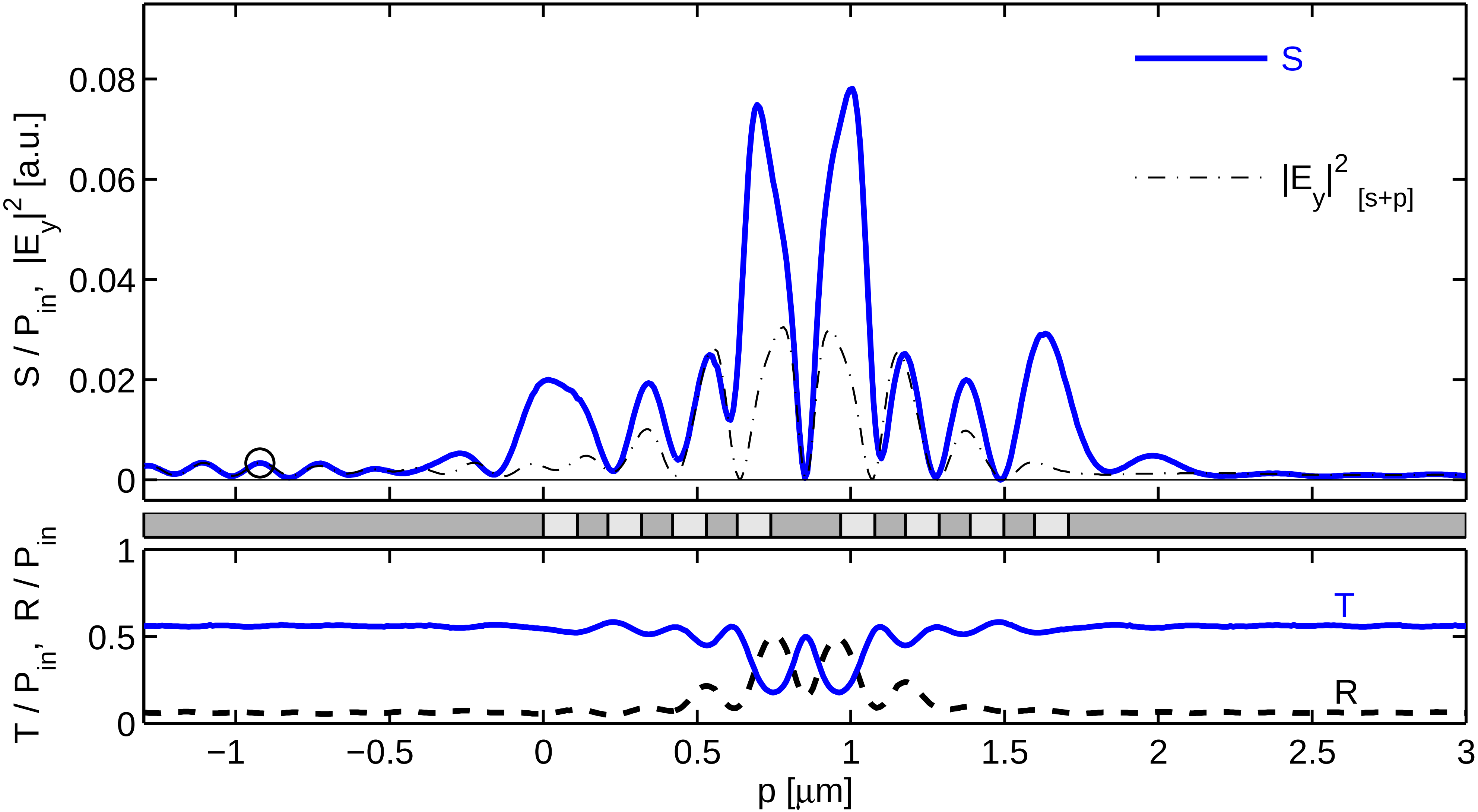


$n_s = 1.45$, $n_f = 2.0$, $n_c = 1.0$, $t = 0.2\,\mu$m, $d = 0.6\,\mu$m,
$\Lambda = 0.21\,\mu$m, $s = 0.11\,\mu$m, $L = 0.2275\,\mu$m, $g = 10\,$nm, $w = 100\,$nm, $n_p = 1.5$,
TE, $\lambda = 0.633\,\mu$m, $(x, z) \in [-3.0, 3.0] \times [-3.0, 4.7075]\,\mu\text{m}^2$, $M_x = 100$, $M_z = 120$.

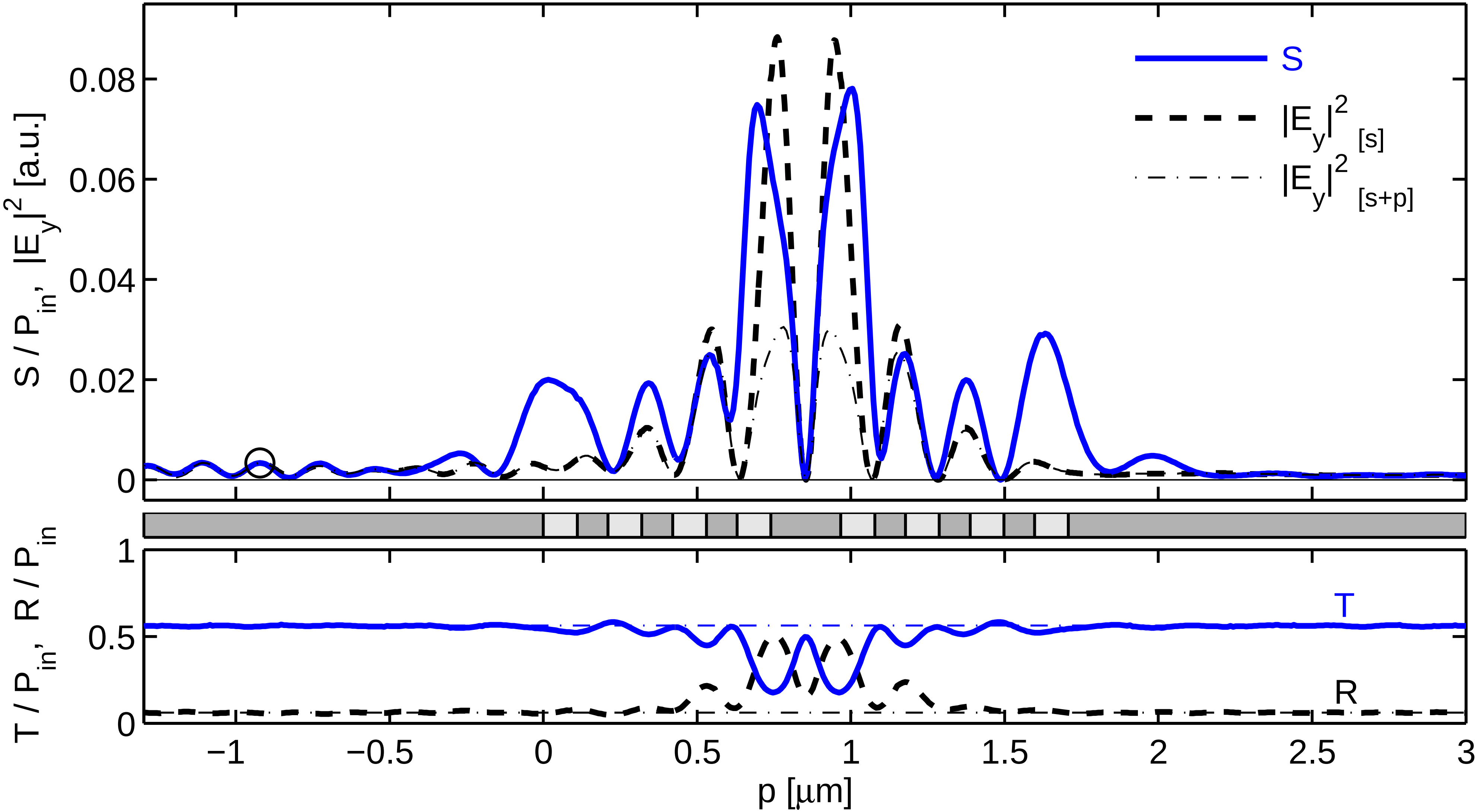
Resonant defect cavity



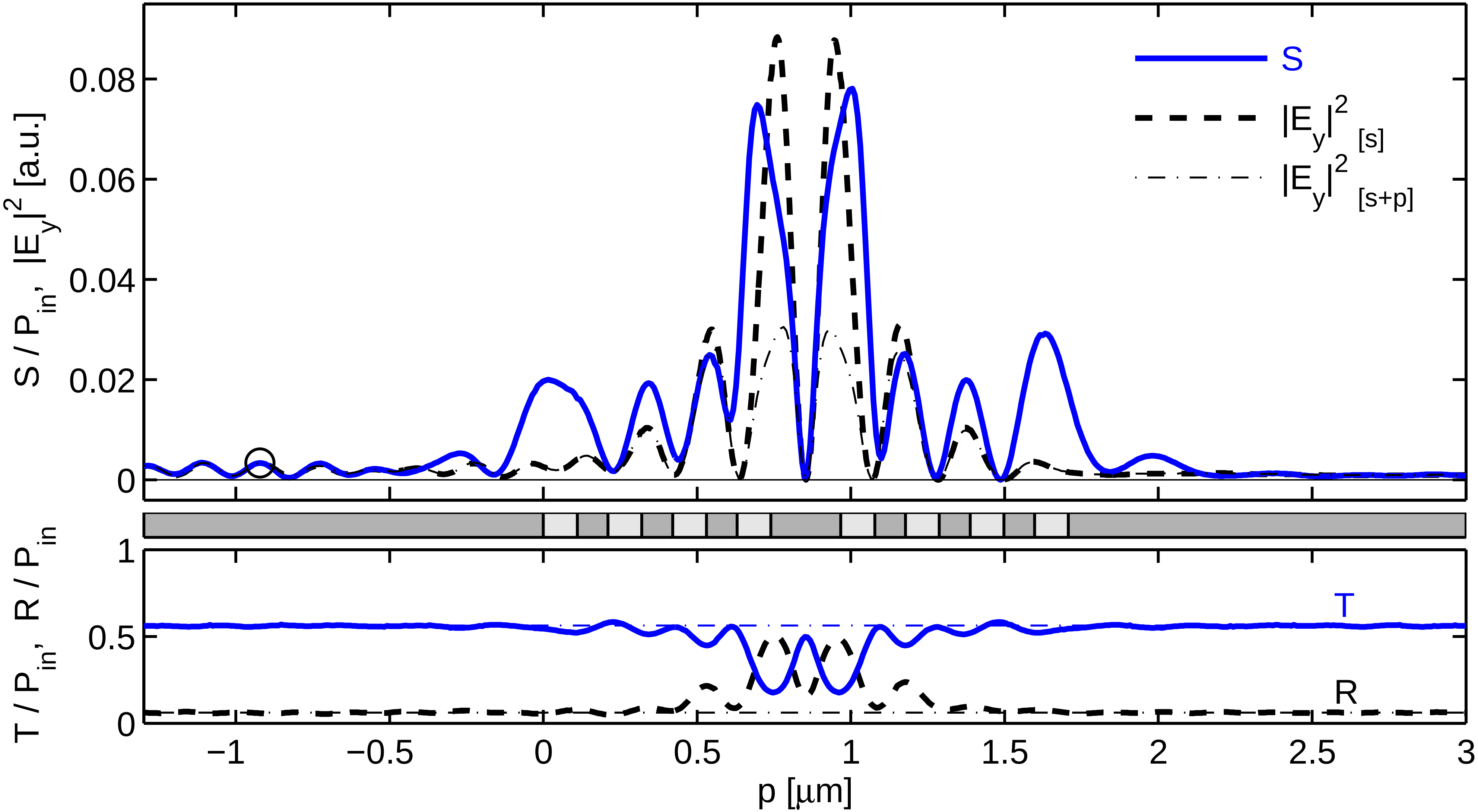
Resonant defect cavity



Resonant defect cavity



Resonant defect cavity

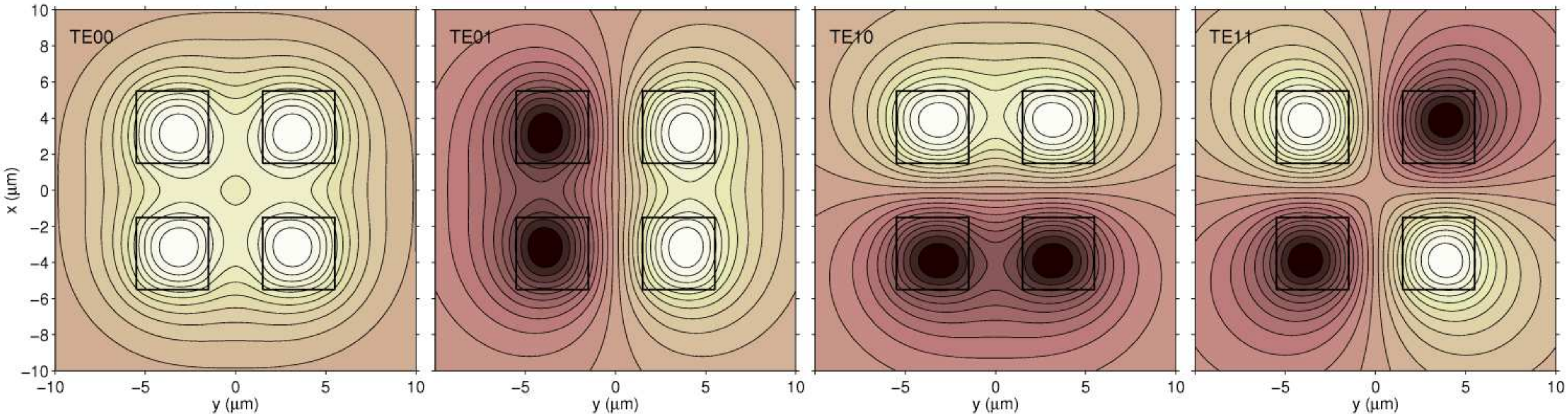


$T(p)$, $R(p)$: the probe switches the transmission.

Outline

- Simulations in guided wave optics
  - Macroscopic Maxwell equations
  - Stationary and time-domain problems
  - 2-D problems
  - Modes of dielectric waveguides
- Quadridirectional eigenmode propagation
  - Problem setting
  - Eigenmode expansion
  - Algebraic procedure
- 2-D PSTM model
  - Probing evanescent fields
  - Hole defect in a slab waveguide
  - Waveguide Bragg grating: experiment & 2-D model
  - Resonant defect cavity
- Optical switching by NEMS-actuated resonator arrays

Ongoing: TOE.7143

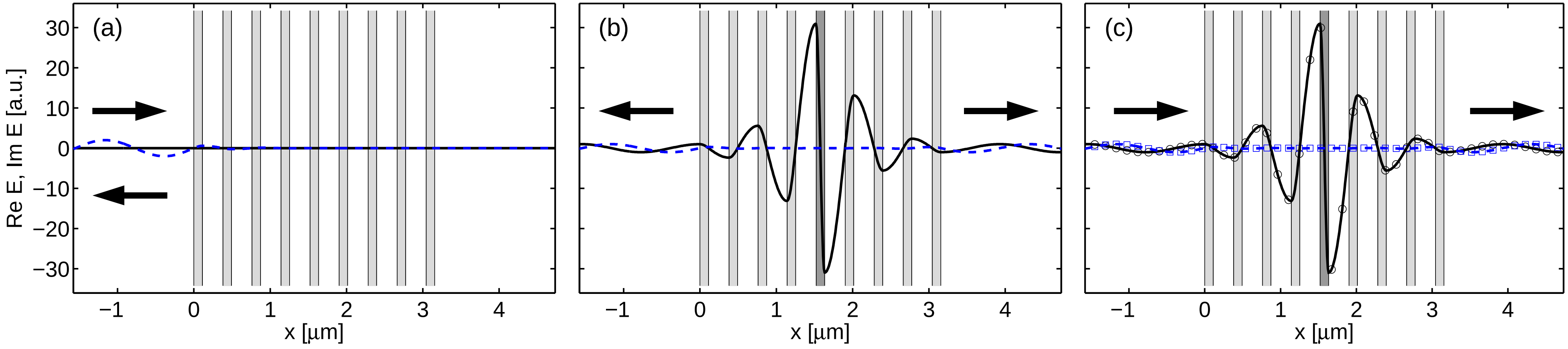


Olena Ivanova

Projection techniques for dimensionality reduction 3-D → 2-D

NanoPhotonics Flagship, project TOE.7143
Optical switching by NEMS actuated resonator arrays — modeling & simulation tools

Ongoing: TOE.7143



Milan Maksimovic

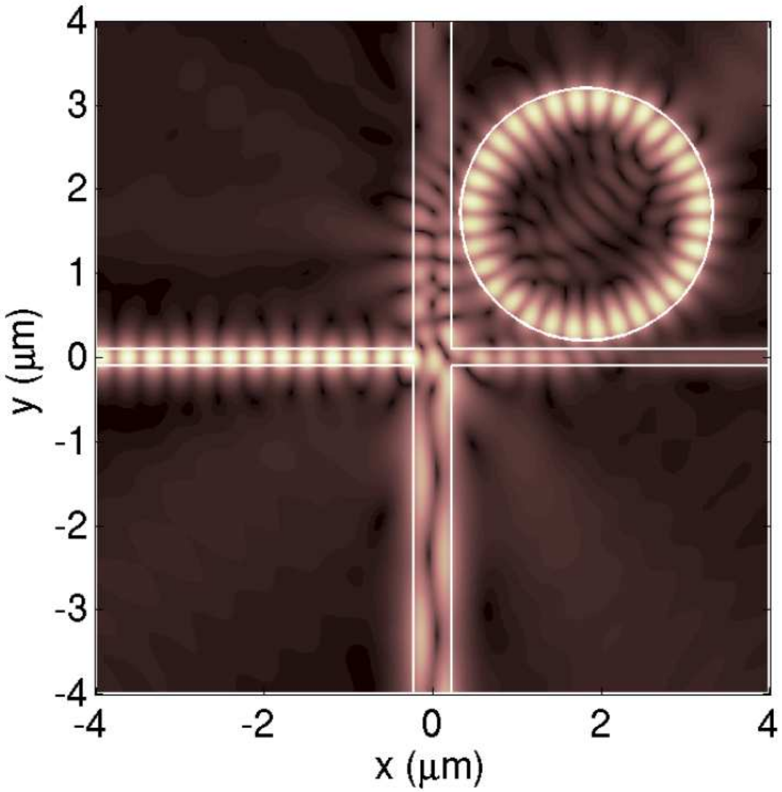
Time domain characterization of resonances & perturbations

nano ned

NanoPhotonics Flagship, project TOE.7143

Optical switching by NEMS actuated resonator arrays — modeling & simulation tools

Ongoing: TOE.7143



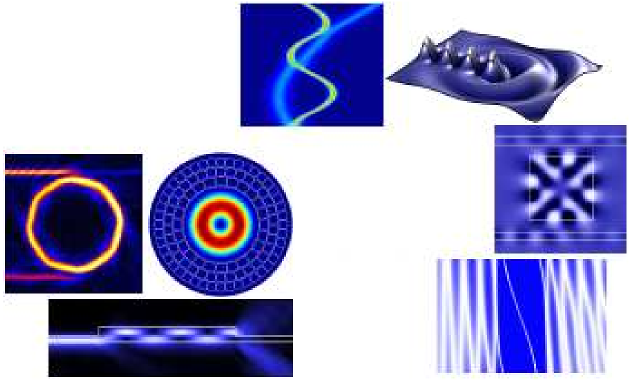
Remco Stoffer

Numerical simulation tools
& interaction with projects
TOE.7144, NEMS, group TST (EEMCS, EL),
TOE.7145, Optics, group IOMS (EEMCS, EL).

nano ned

NanoPhotonics Flagship, project TOE.7143
Optical switching by NEMS actuated resonator arrays — modeling & simulation tools

Photonics @ AAMP



http://www.math.utwente.nl/aamp/

http://www.math.utwente.nl/aamp/research.html#MATHOPT